

# Modern Physics

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Lesson VI  
September 28, 2023

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- 1 Particle Dynamics
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  - Conservation of 4-momentum and all that...
  - Two-body decay of unstable particles



- Newton's first law of motion holds in special relativistic mechanics as well as nonrelativistic mechanics
- In absence of forces  $\Rightarrow$  body is at rest or moves in straight line at constant speed

$$\frac{d\mathbf{U}}{d\tau} = 0 \quad (1)$$

$\mathbf{U}(\gamma c, \gamma \vec{u}) \Rightarrow$  (1) implies that  $\vec{u}$  is constant in any inertial frame

- Objective of relativistic mechanics:  $\Rightarrow$  introduce the analog of Newton's 2nd law

$$\vec{F} = m\vec{a} \quad (2)$$

- There is nothing from which this law can be derived but plausibly it must satisfy certain properties:
  - 1 It must satisfy the principle of relativity  
i.e.  $\Rightarrow$  take the same form in every inertial frame
  - 2 It must reduce to (1) when the force is zero
  - 3 It must reduce to (2) in any inertial frame  
when speed of particle is much less than speed of light

$$m \frac{d\mathbf{U}}{d\tau} = \mathbf{f} \quad (3)$$

- $m$  characterizes particle's inertial properties
- $\mathbf{f}$  4-force.
- Using  $d\mathbf{U}/d\tau = \mathbf{A}$  (3) can be rewritten in evocative form

$$\mathbf{f} = m\mathbf{A} \quad (4)$$

- This represents 4-equations but they are not all independent
- Normalization of the 4-velocity  $U_\mu U^\mu = c^2$  implies

$$m \frac{d(\mathbf{U} \cdot \mathbf{U})}{d\tau} = 0 \Rightarrow \mathbf{f} \cdot \mathbf{U} = 0 \quad (5)$$

- (5) shows there are only 3 independent equations of motion  
– same number as in Newtonian mechanics –

4-momentum defined by

$$\mathbf{p} = m \mathbf{U} \quad (6)$$

equation of motion can be rewritten as

$$\frac{d\mathbf{p}}{d\tau} = \mathbf{f} \quad (7)$$

important property of 4-momentum  $\Rightarrow$  invariant mass

$$p_\mu p^\mu = m^2 c^2 \quad (8)$$

- Components of 4-momentum related  $\vec{u}$  according to

$$p^0 = \frac{mc}{\sqrt{1 - u^2/c^2}} \quad \text{and} \quad \vec{p} = \frac{m\vec{u}}{\sqrt{1 - u^2/c^2}} \quad (9)$$

- For small speeds  $u \ll c$

$$p^0 = mc + \frac{1}{2}m\frac{u^2}{c} + \dots \quad \text{and} \quad \vec{p} = m\vec{u} + \dots \quad (10)$$

- $\vec{p}$  reduces to usual 3-momentum
- $p^0$  reduces to kinetic energy per units of  $c$  plus mass in units of  $c$
- $p^\mu$  also called energy-momentum 4-vector

$$p^\mu = (E/c, \vec{p}) = (m\gamma c, m\gamma\vec{u}) \quad (11)$$

- mass is part of energy of relativistic particle

$$p_\mu p^\mu = m^2 c^2 \Rightarrow E = (m^2 c^4 + \vec{p}^2 c^2)^{1/2} \quad (12)$$

- For particle at rest (12) reduces to  $E = mc^2$

In particular inertial frame...

- Connection between relativistic equation of motion and Newton's laws can be made more explicit by defining 3-force

$$\frac{d\vec{p}}{dt} \equiv \vec{F} \quad (13)$$

- Same form as Newton's law but with relativistic expression for  $\vec{p}$
- Only difference  $\Rightarrow$  different relation momentum to velocity

$$\vec{f} = \frac{d\vec{p}}{d\tau} = \frac{d\vec{p}/dt}{dt/d\tau} = \gamma\vec{F} \quad (14)$$

- 4-force acting on particle can be written in terms of 3-force

$$f = (\gamma\vec{F} \cdot \vec{u}, \gamma\vec{F}) \quad (15)$$

- Time component of equation of motion

$$\frac{dE}{dt} = \vec{F} \cdot \vec{u} \quad (16)$$

familiar relation from Newtonian mechanics



- Time component of equation of motion consequence of other 3
- In terms of three force  $\Rightarrow$  equations of motion take same form as they do in usual Newtonian mechanics but with relativistic expressions for energy and momentum
- For  $v \ll c \Rightarrow$  relativistic version of Newton's second law reduces to the familiar nonrelativistic form
- Newtonian mechanics is low-velocity approximation of relativistic mechanics

## Light rays

- Massless particles move at speed of light along null trajectories
- Proper time interval between any two points  $\Rightarrow$  is zero
- Curve  $x = ct$  could be written parametrically in term of  $\lambda$

$$x^\mu = U^\mu \lambda \quad \text{with} \quad U^\mu = (c, c, 0, 0) \quad (17)$$

- $U$  is a null vector  $\Rightarrow$

$$U \cdot U = 0 \quad (18)$$

- Different parametrizations give different tangent 4-vectors  
but all have zero length
- With choice (17)

$$\frac{dU}{d\lambda} = 0 \quad (19)$$

light ray equation of motion is same as for particle

- Basic law of collision mechanics  $\Rightarrow$  conservation of 4-momentum:  
*Sum of 4-momenta of all particles going into point-collision  
 is same as sum of 4-momenta of all those coming out*

$$\sum^* p_i = 0 \quad (20)$$

$\sum^*$   $\Rightarrow$  sum that counts pre-collision terms positively  
 and post-collision terms negatively

- For closed system  $\Rightarrow$  conservation of total 4-momentum  
 can be shown to be result of spacetime homogeneity
- Whether law is actually true must be decided by experiment
- Countless experiments have shown  
 that total 4-momentum of isolated system is constant

## CM frame

- If we have a system of particles with 4-momenta  $p_i$   
subject to no forces except mutual collisions
- total 4-momentum  $p_{\text{tot}} = \sum p_i$  is timelike and future-pointing
- $\Rightarrow$  there exists an inertial frame  $S$   
in which spatial components of  $p_{\text{tot}}$  vanish

$S$  should be called center-of-momentum frame but is called CM frame

## Invariant mass

- Invariant mass of two particles with 4-momenta  $p_a$  and  $p_b$

$$m_{ab}^2 c^2 = (p_a + p_b)^2$$

- Invariant mass useful to find mass of short-live unstable particles from momenta of their observed decay products
- Consider  $X \rightarrow a + b \Leftrightarrow p_X = p_a + p_b$

$$\begin{aligned} m_X^2 c^2 &= (p_a + p_b)^2 = p_a^2 + p_b^2 + 2p_a \cdot p_b \\ &= m_a^2 c^2 + m_b^2 c^2 + 2E_a E_b / c^2 - 2\vec{p}_a \cdot \vec{p}_b \end{aligned} \quad (21)$$

- In high energy experiment  
3-momenta and masses of particles  $a$  and  $b$  must be measured
- For charged particles this requires a magnetic field  
and tracking of trajectory to measure bending  
as well as some means of particle identification
- One must also identify the vertex and measure the opening angle

- Consider decay process  $X \rightarrow ab$
- In CM frame for  $a$  and  $b$  mother particle  $X$  is at rest
- 4-momenta

$$\mathbf{p}_X = (Mc, 0, 0, 0) \quad \mathbf{p}_a = (E_a/c, \vec{p}_a) \quad \mathbf{p}_b = (E_b/c, \vec{p}_b) \quad (22)$$

- Conservation of 4-momentum requires:

$$\mathbf{p}_X = \mathbf{p}_a + \mathbf{p}_b \quad \Rightarrow \quad \vec{p}_a = -\vec{p}_b$$

- Omitting subscript on 4-momenta  $\Rightarrow$  energy conservation reads

$$E_a + E_b = \sqrt{m_a^2 c^4 + p^2 c^2} + \sqrt{m_b^2 c^4 + p^2 c^2} = Mc^2 \quad (23)$$

- Solving (23) for  $p$

$$p = c \frac{\sqrt{[M^2 - (m_a - m_b)^2][M^2 - (m_a + m_b)^2]}}{2M} \quad (24)$$

- Immediate consequence

$$M \geq m_a + m_b \quad (25)$$

$X$  decays only if mass exceeds sum of decay products masses

- Conversely
  - ☞ if particle has mass exceeding masses of two other particles  
**particle is unstable and decays**  
 unless decay is forbidden by some conservation law  
 e.g. conservation of charge, momentum, and angular momentum
- Momenta of daughter particles and energies fixed by 3 masses:  
 from energy conservation (23) ☞  $E_b = \sqrt{E_a^2 - m_a^2 c^4 + m_b^2 c^4}$   
 solve to get

$$E_a = \frac{1}{2M} (M^2 + m_a^2 - m_b^2) c^2 \quad (26)$$

similarly

$$E_b = \frac{1}{2M} (M^2 + m_b^2 - m_a^2) c^2 \quad (27)$$

- No preferred direction in which the daughter particles travel  
**decay is said to be *isotropic***  
 ☞ daughter particles travelling *back-to-back* in  $X$  rest frame

- Of interest is also two-body decay of unstable particles in flight
- In-flight decays  $\Rightarrow$  only way to measure mass of neutral particle
- Take  $z$ -axis along direction of flight of mother particle

$$\mathbf{p}_X = (E/c, 0, 0, p) \quad \mathbf{p}_a = (E_a/c, \vec{p}_{a\perp}, p_{az}) \quad \mathbf{p}_b = (E_b/c, \vec{p}_{b\perp}, p_{bz})$$

- By momentum conservation  $\Rightarrow$  transverse momentum vectors

$$\vec{p}_{\perp} \equiv \vec{p}_{a\perp} = -\vec{p}_{b\perp} \quad (28)$$

- Energies and  $z$  components of particle momenta related to those in the CM frame by a Lorentz boost with a boost velocity equal to the speed of the mother particle

$$E_a/c = \gamma(E_a^*/c + \beta p_{az}^*)$$

$$p_{az} = \gamma(p_{az}^* + \beta E_a^*/c)$$

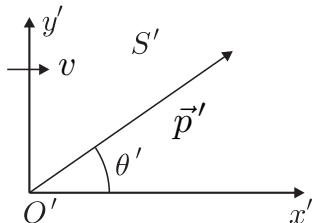
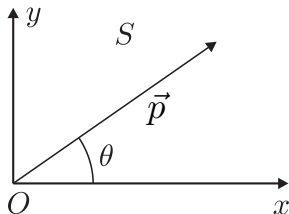
$$\vec{p}_{a\perp} = \vec{p}_{a\perp}^*$$

$$\beta = pc/E \text{ and } \gamma = E/(Mc^2)$$

similarly for particle  $b$



- This completely solves problem  
e.g.  $\Rightarrow$  we can find angles which daughter particles  
make with  $z$ -axis and with each other as functions of  $p_x$



- In  $S$  we have

$$p^\mu = (E/c, p \cos \theta, p \sin \theta, 0) \quad (29)$$

in  $S'$  it follows that

$$p'^\mu = (E'/c, p' \cos \theta', p' \sin \theta', 0) \quad (30)$$

Applying  $S \rightarrow S'$  Lorentz transformation

$$\begin{aligned} p' \cos \theta' &= \gamma^* (p \cos \theta - \beta^* E/c) \\ p' \sin \theta' &= p \sin \theta \end{aligned} \quad (31)$$

so

$$\tan \theta' = \frac{p \sin \theta}{\gamma^* (p \cos \theta - \beta^* E/c)} \quad (32)$$

or

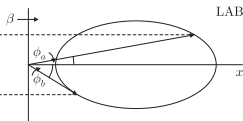
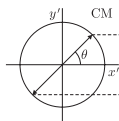
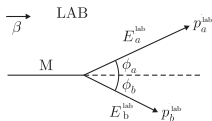
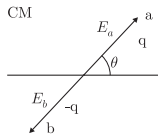
$$\tan \theta' = \frac{\sin \theta}{\gamma^* (\cos \theta - \beta^*/\beta)} \quad (33)$$

$\beta^* = v/c$   $\Rightarrow$  velocity of  $S'$  wrt  $S$  and  $\beta = pc/E$   $\Rightarrow$  velocity of particle in  $S$

Inverse relation is found to be

$$\tan \theta = \frac{\sin \theta'}{\gamma^* (\cos \theta' + \beta^*/\beta')} \quad (34)$$

$\beta' = p'c/E'$   $\Rightarrow$  velocity of particle in  $S'$



## Example

- Mother particle of mass  $M$  is traveling with velocity  $\beta = |pc|/E$
- In the CM frame  $\Rightarrow$  particle  $a$  has energy  $E_a$  and momentum  $\vec{q}$  @ angle  $\theta$  with respect to  $x'$ -axis
- Particle momenta in the lab frame

$$\begin{aligned} p_a^{\text{lab}} \cos \phi_a &= \gamma(q \cos \theta + \beta E_a/c) \\ p_b^{\text{lab}} \cos \phi_b &= \gamma(-q \cos \theta + \beta E_b/c) \end{aligned} \quad (35)$$

- Use inverse Lorentz transformation to obtain variables in CM frame from measured parameters in lab

- 2nd approach  $\Rightarrow$  start from energy-momentum conservation

$$E = E_a + E_b = \sqrt{m_a^2 c^4 + p_a^2 c^2} + \sqrt{m_b^2 c^4 + p_b^2 c^2} \quad (36)$$

$$\vec{p} = \vec{p}_a + \vec{p}_b \quad (37)$$

- Substituting in (36)  $p_b^2$  by  $(\vec{p} - \vec{p}_a)^2$

$$p_a = \frac{(M^2 + m_a^2 - m_b^2)c^2 p \cos \theta_a \pm 2E \sqrt{M^2 p^{*2} - m_a^2 p^2 \sin^2 \theta_a}}{2(M^2 c^2 + p^2 \sin^2 \theta_a)}$$

- By demanding  $p_a$  to be real  $\Rightarrow M^2 p^{*2} - m_a^2 p^2 \sin^2 \theta_a \geq 0$
- This condition is satisfied for all angles  $\theta_a$  if  $Mp^*/(m_a p) > 1$   
negative sign must be rejected  $\Rightarrow$  unphysical  $p_a < 0$  for  $\theta_a > \pi/2$
- If  $Mp^*/(m_a p) < 1$   
 $\Rightarrow$  region of parameter space in which both signs must be kept:  
for each value of  $\theta_a < \theta_{a,\max}$  there are two values of  $p_a$   
and correspondingly also two values of  $p_b$