



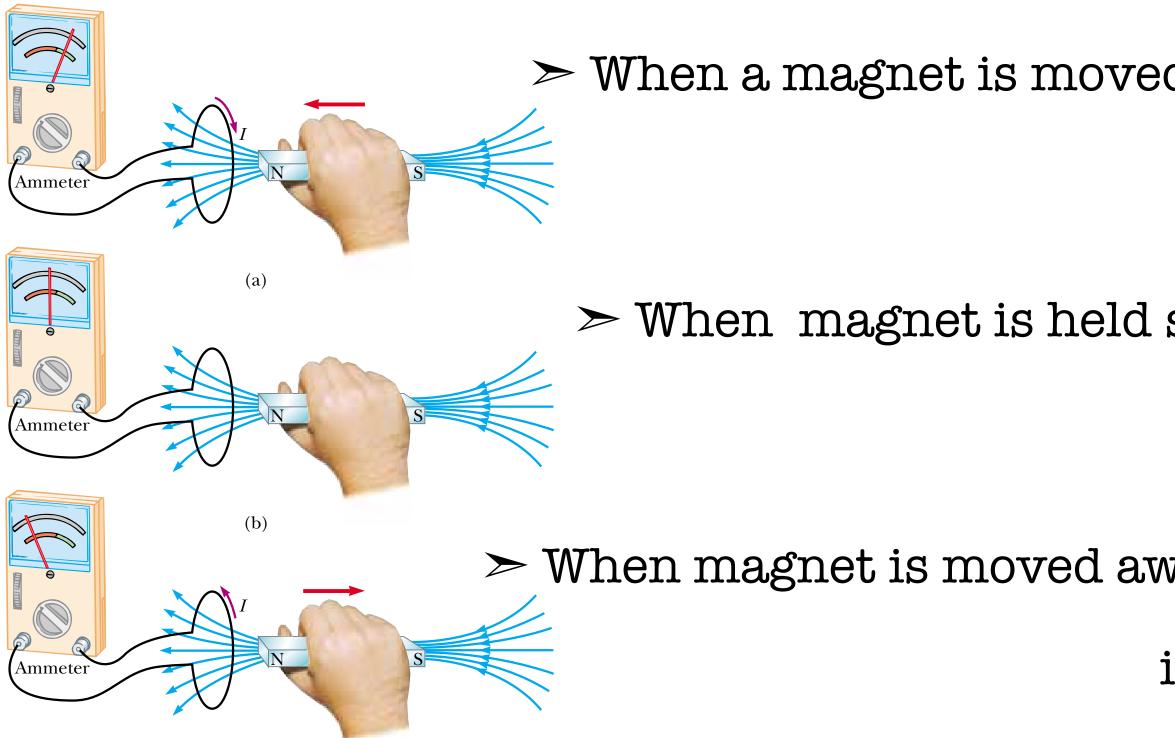
Question 🖛

> Steady electric current can give steady magnetic field

> Because of symmetry between electricity and magnetism - we can ask:

> Steady magnetic field can give steady electric current

> OR Changing magnetic field can give steady electric current



- \succ When a magnet is moved toward a loop of wire \leftarrow sensitive ammeter deflects indicating that current is induced in the loop
 - > When magnet is held stationary $rac{r}{r}$ there is no induced current in the loop even when the magnet is inside the loop
- > When magnet is moved away from loop $rac{r}$ ammeter deflects in opposite direction indicating that induced current is in opposite direction



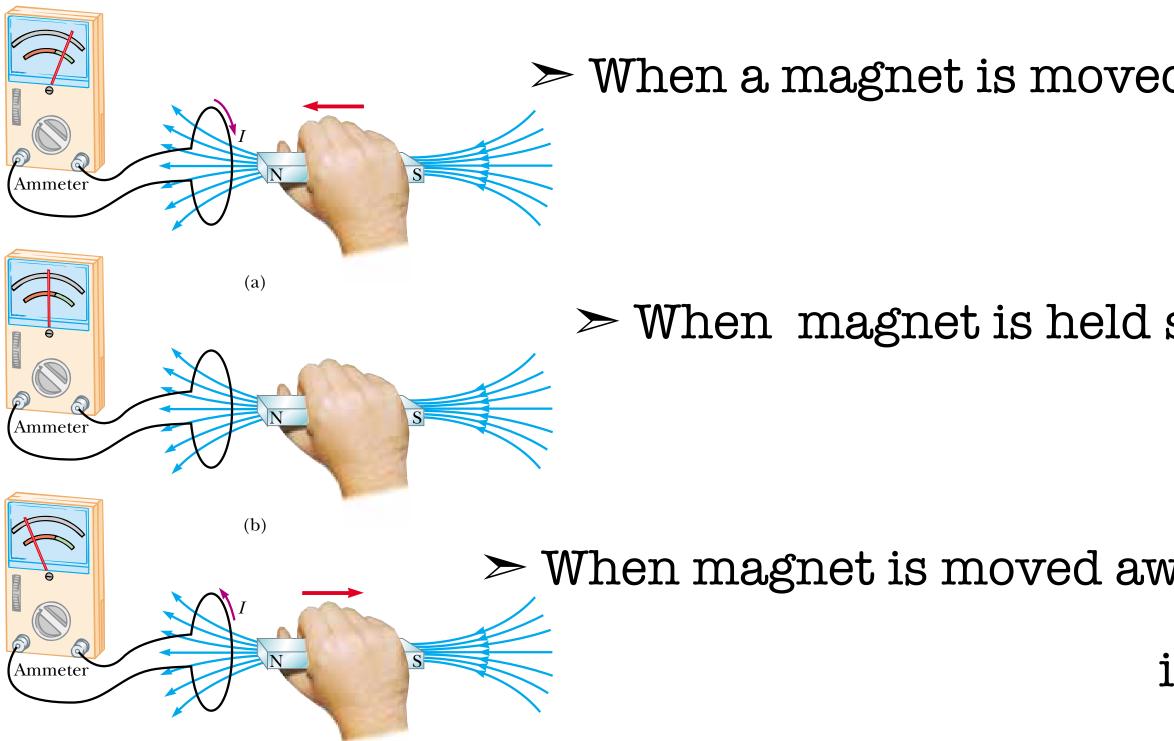
ANSWER 🖛

> Steady electric current can give steady magnetic field

> Because of symmetry between electricity and magnetism - we can ask:

> Steady magnetic field can give steady electric current X

> OR Changing magnetic field can give steady electric current \checkmark

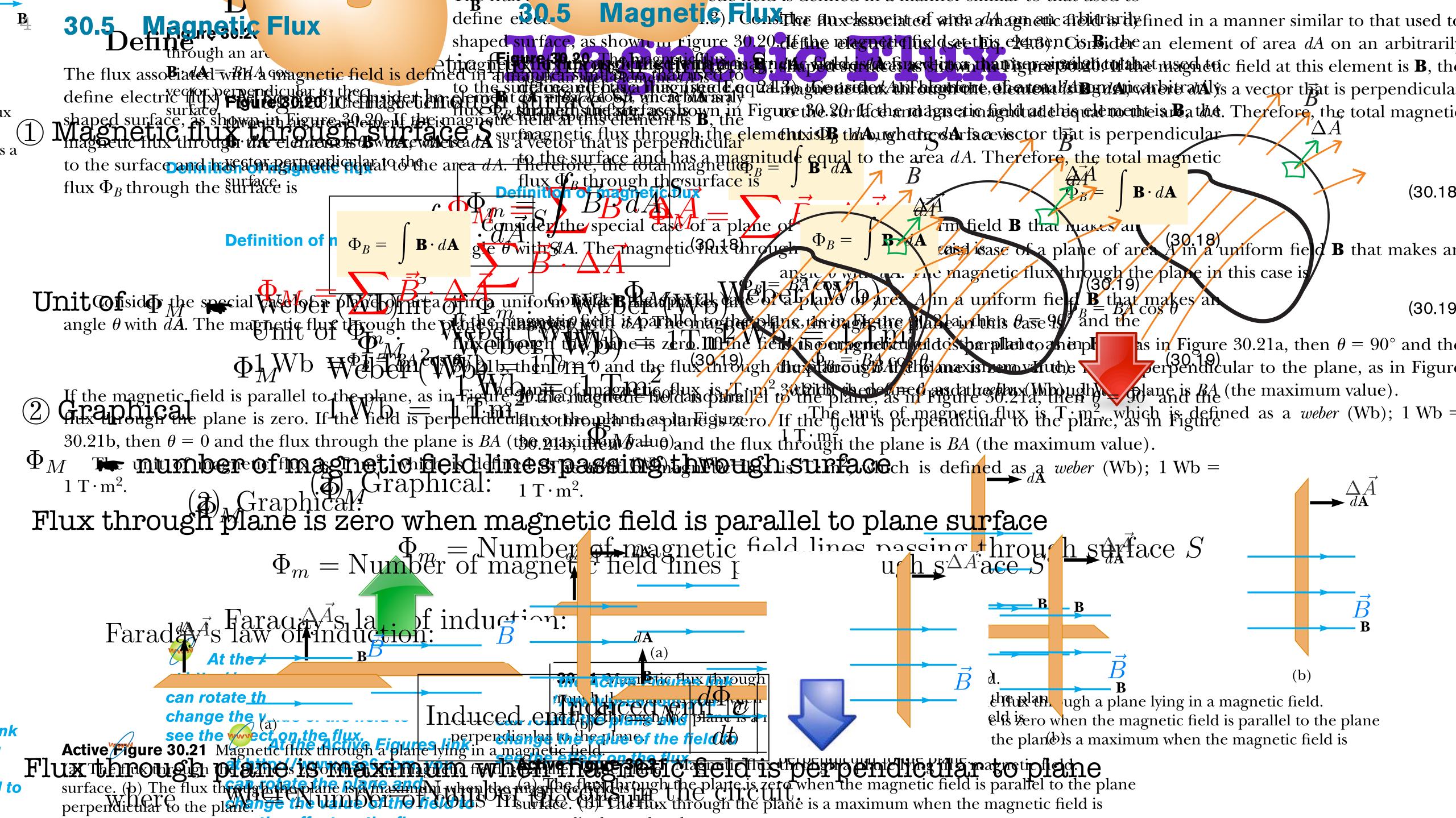


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> The net magn

The net magnetic flux Φ_B through any closed surface is equal to zero:

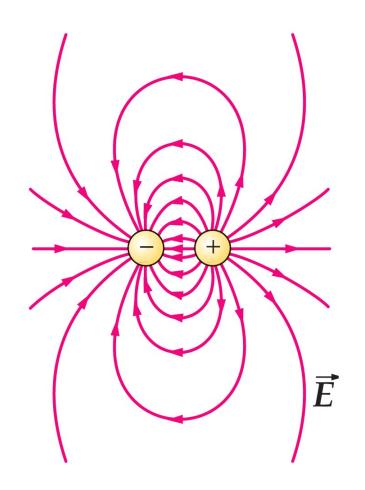
Gauss's Law for Magnetic Field

The net magnetic flux Φ_B through any closed surface is equal to zero:

- \succ There a
- ≻ Magnet
- ≻ No matt

There are no magnetic charges. Magnetic field lines always close in themselves. No matter how the (closed) Gaussian surface is chosen, the net magnetic flux through it always vanishes.

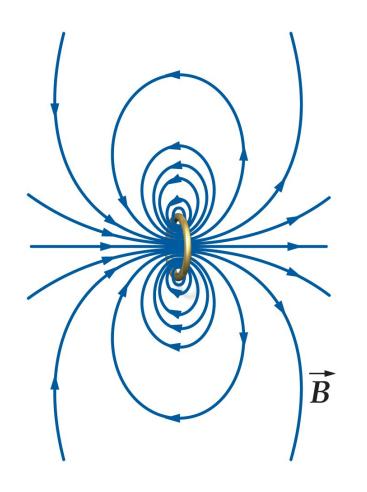
The figures below illustrate Gauss's laws for the electric and magnetic fields in the context of an electric dipole (left) and a magnetic dipole (right).

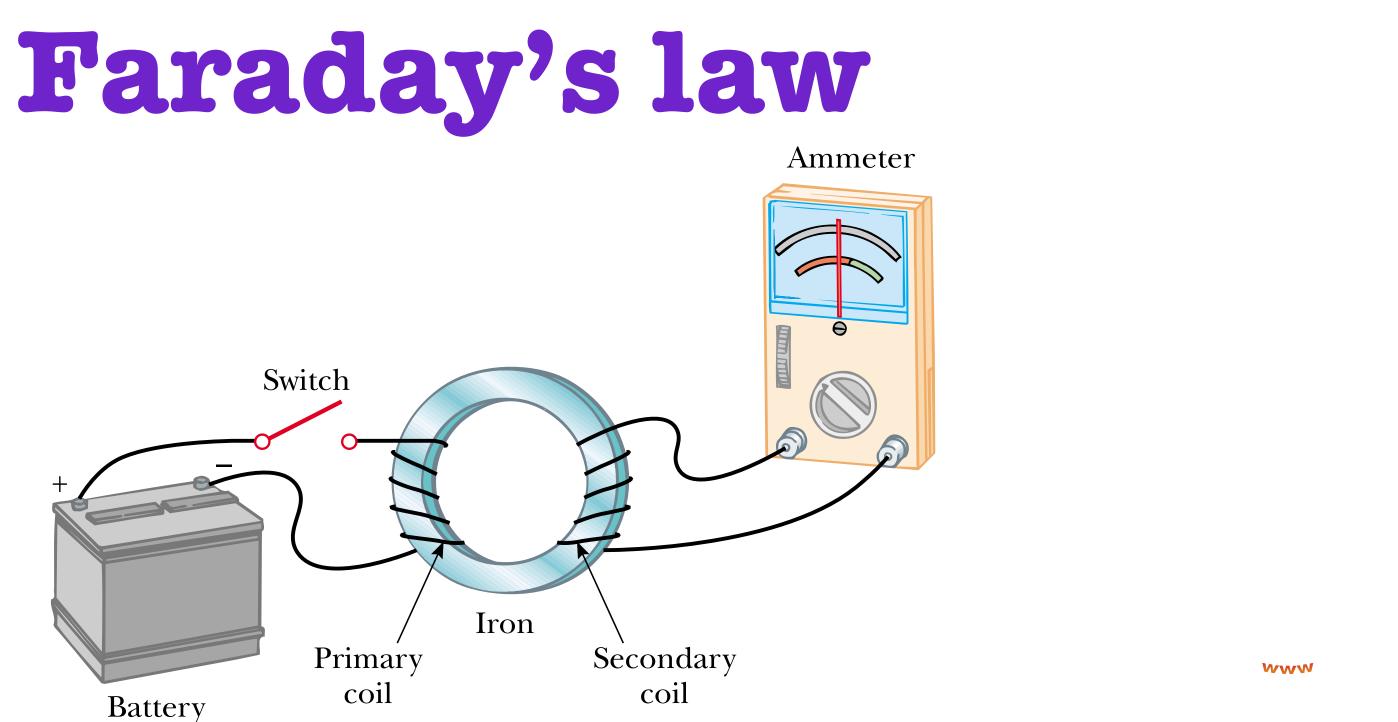


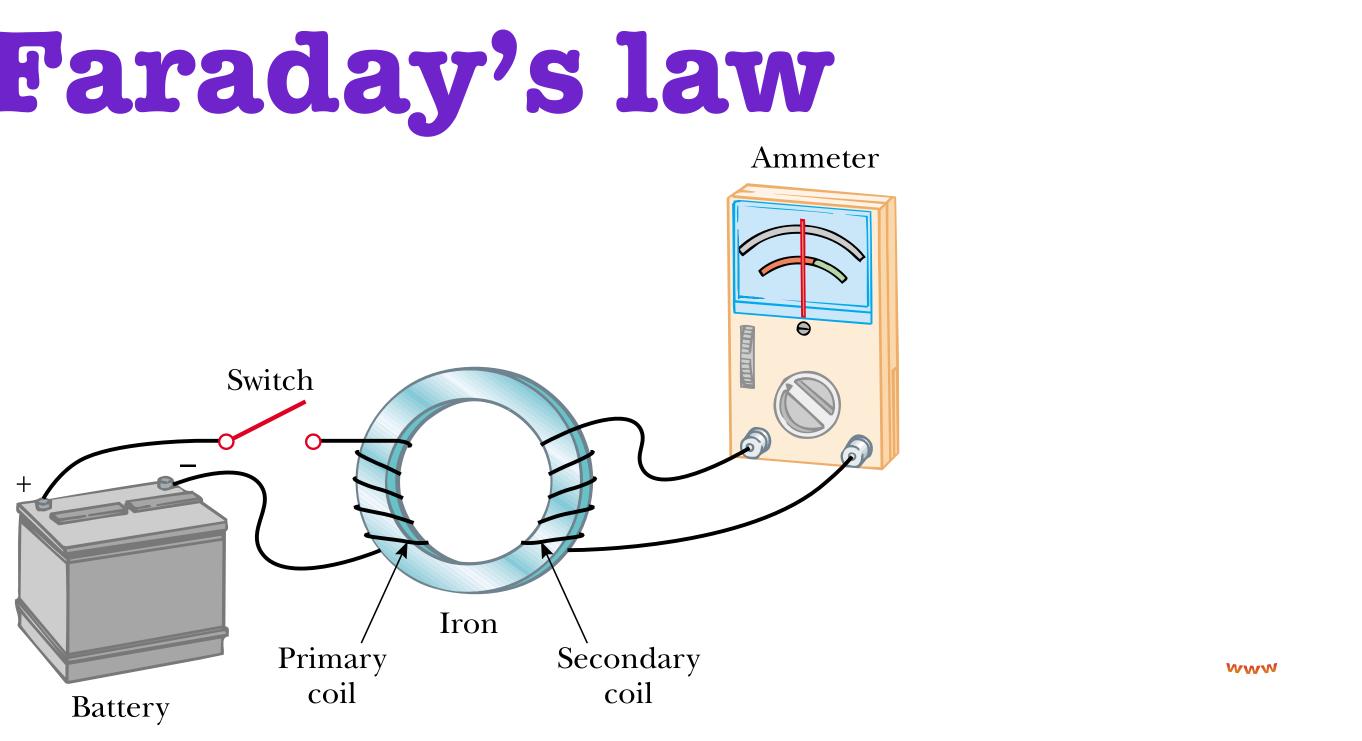




$$\oint \vec{B} \cdot d\vec{A} = 0.$$







Faraday's experiment 📫

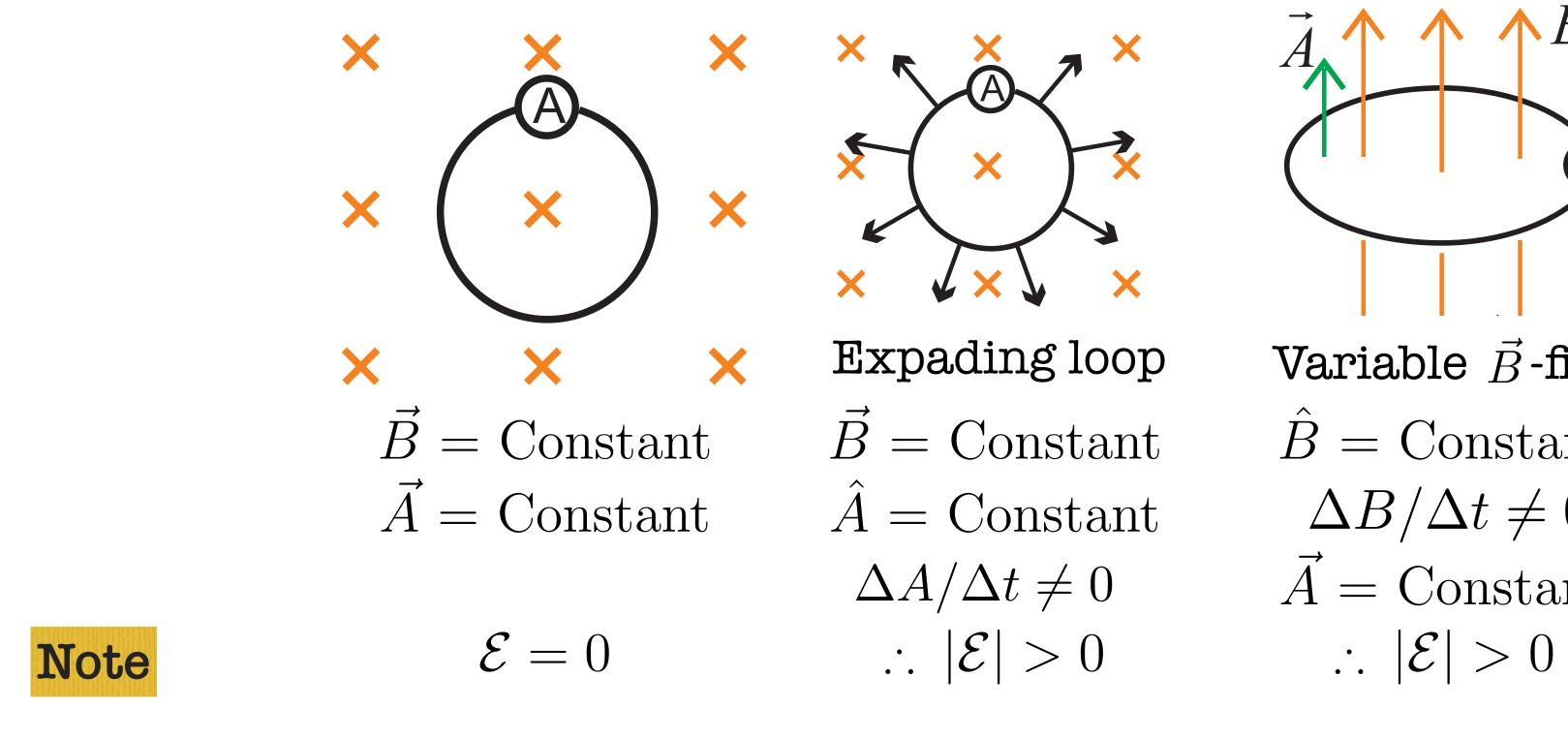
When switch in primary circuit is closed ammeter in secondary circuit deflects momentarily

emf induced in secondary circuit is caused by changing magnetic field through secondary coil

Faraday's law of induction \Rightarrow Induced emf \Rightarrow \mathcal{E}

number of coils in circuit

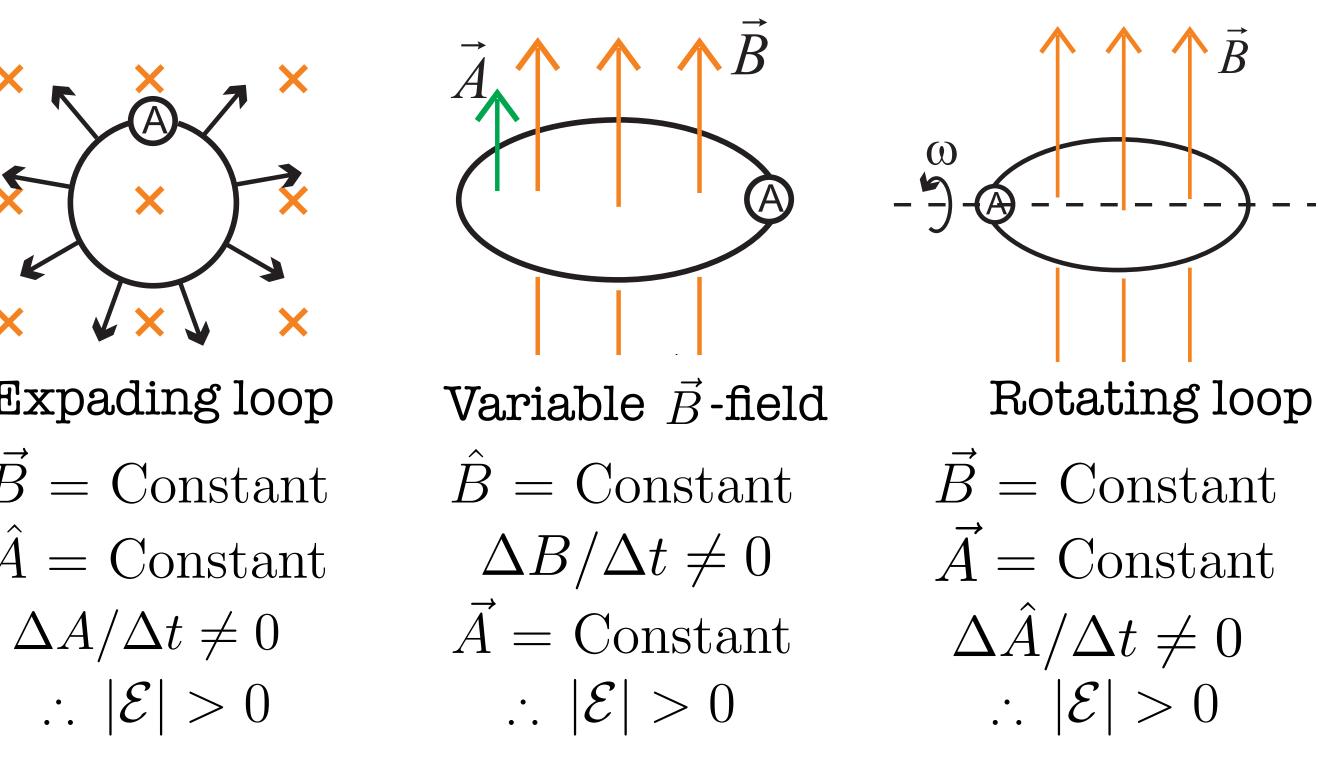
$$| = N \left| \frac{\Delta \Phi_M}{\Delta t} \right| \qquad |\mathcal{E}| = N \left| \frac{\Delta \Phi_M}{\Delta t} \right|$$



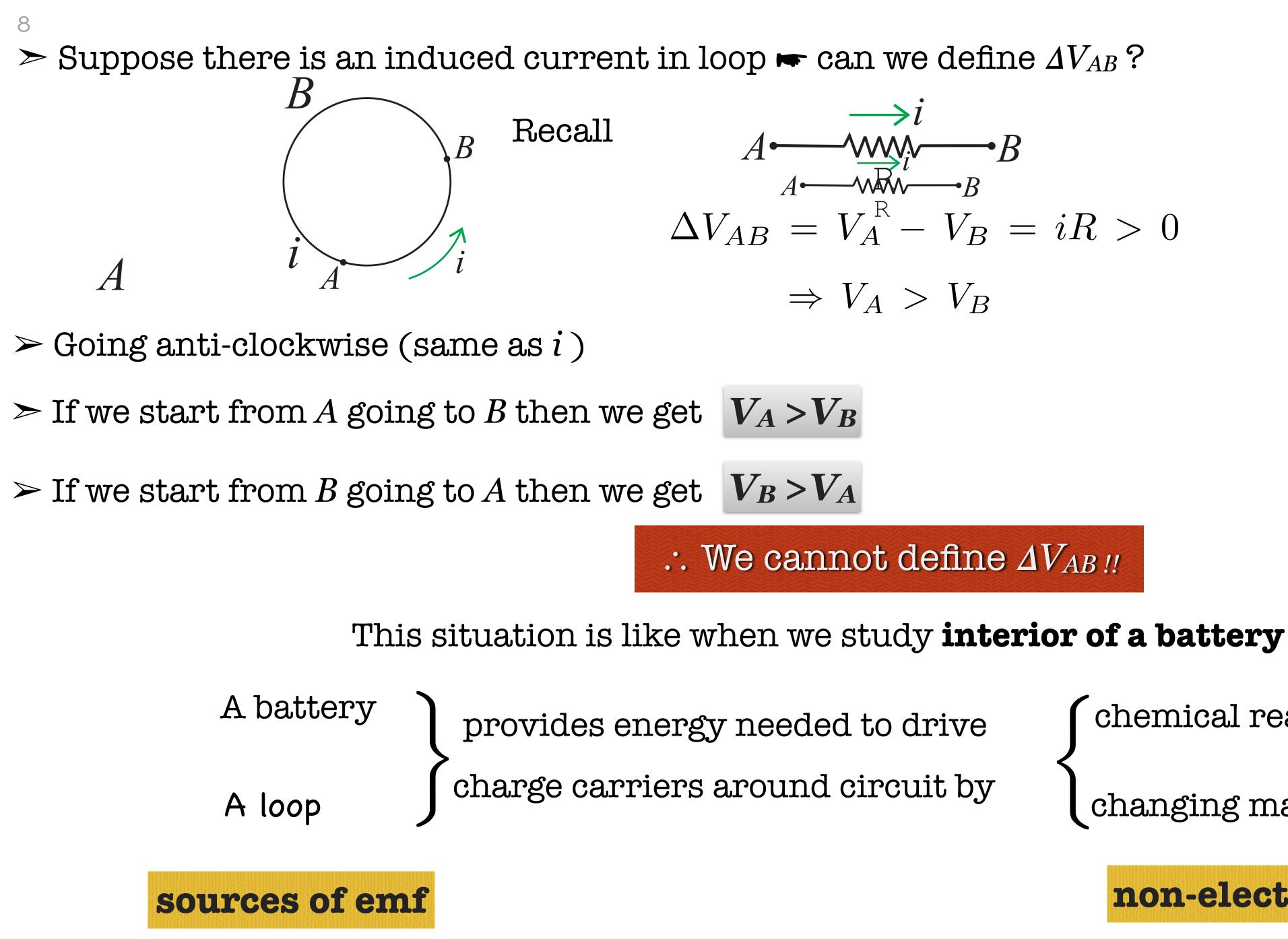
> Induced emf drives a current throughout circuit similar to function of a battery

> Difference here is that induced emf is distributed throughout circuit consequence

we cannot define a potential difference between any two points in circuit







chemical reactions

changing magnetic flux

non-electric means



(1) Flux of magnetic field due to induced current **opposes** change in flux that causes induced current

(2) Induced current is in such a direction as to **oppose** changes that produces it

(3) Incorporating Lenz's law into Faraday's Law 🖛

If
$$\frac{\Delta \Phi_M}{\Delta t} > 0 \Rightarrow \Phi_M \uparrow \mathcal{E} \text{ appears} \Rightarrow \text{Induc}$$

 $\Rightarrow \vec{B}$ -field due to induced current \Rightarrow calling in



$$\mathcal{E} = -N \frac{\Delta \Phi_M}{\Delta t}$$

ced current appears

so that
$$\Phi_M \Rightarrow \Phi_M \downarrow$$



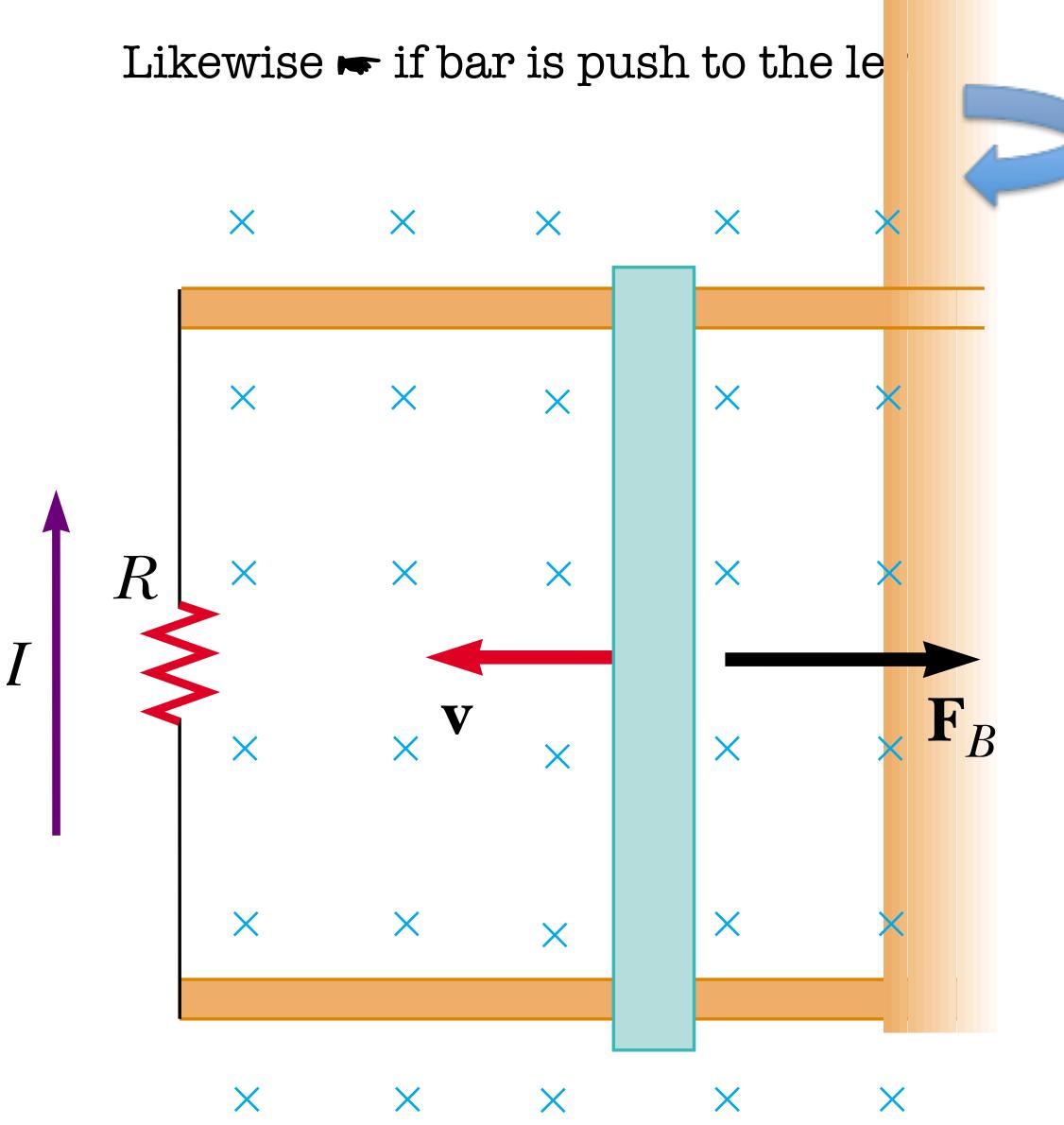
Lenz's Law is consequence from **principle of conservation of energy** (4)

- Suppose bar is given slight push to right
- > This motion sets up a counterclockwise current in the loop

- > What happens if we assume that current is clockwise such that direction of magnetic force exerted on bar is to the right?
- > This force would accelerate the rod and increase its velocity
- > This (in turn) would cause area enclosed by loop to increase more rapidly this would result in increase in induced current which would cause increase in force which would produce increase in current ... and so on...
- > System would acquire energy with no input of energy
- > This is clearly inconsistent with all experience and violates law of energy conservation
- \succ We are forced to conclude that current must be counterclockwise

- X \times X X \times X \times X \times \times X \times X X

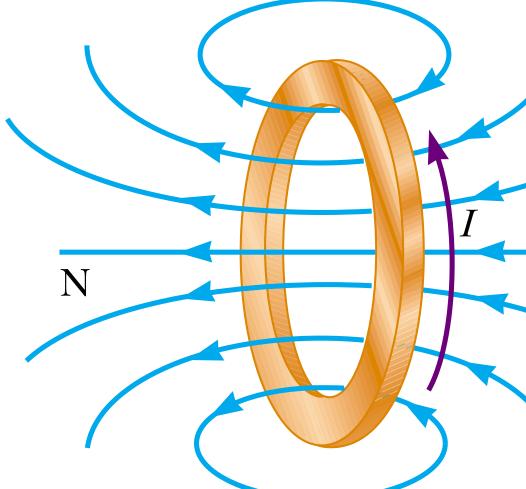
 $\mathbf{B}_{in_{y}}$





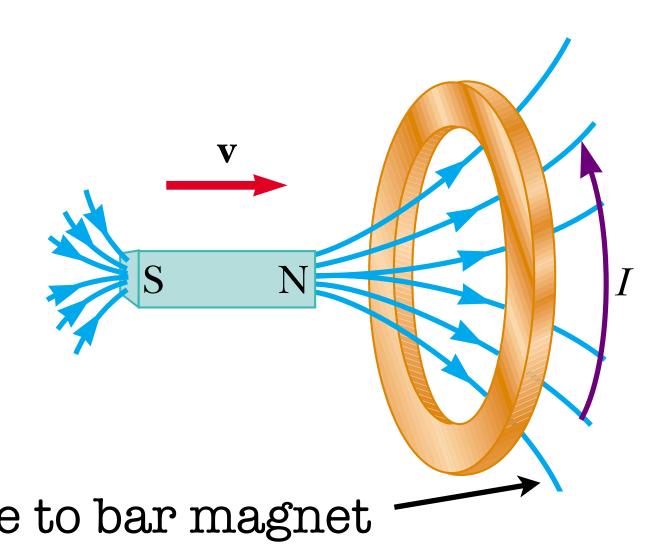
> Magnetic field lines shown are those due to bar magnet

> This induced current produces its own magnetic field directed to the left



> Magnetic field lines shown are those due to induced current in ring

> When magnet is moved toward stationary conducting loop current is induced in the direction shown





- - that counteracts the increasing external flux





e bar moves to the left, the duced current must be clockwise. _{ny?} > When magnet is moved away from stationary conducting loop current is induced in direction shown Ν

 $\overset{(a)}{\succ}$ Magnetic field lines shown are those due to bar magnet

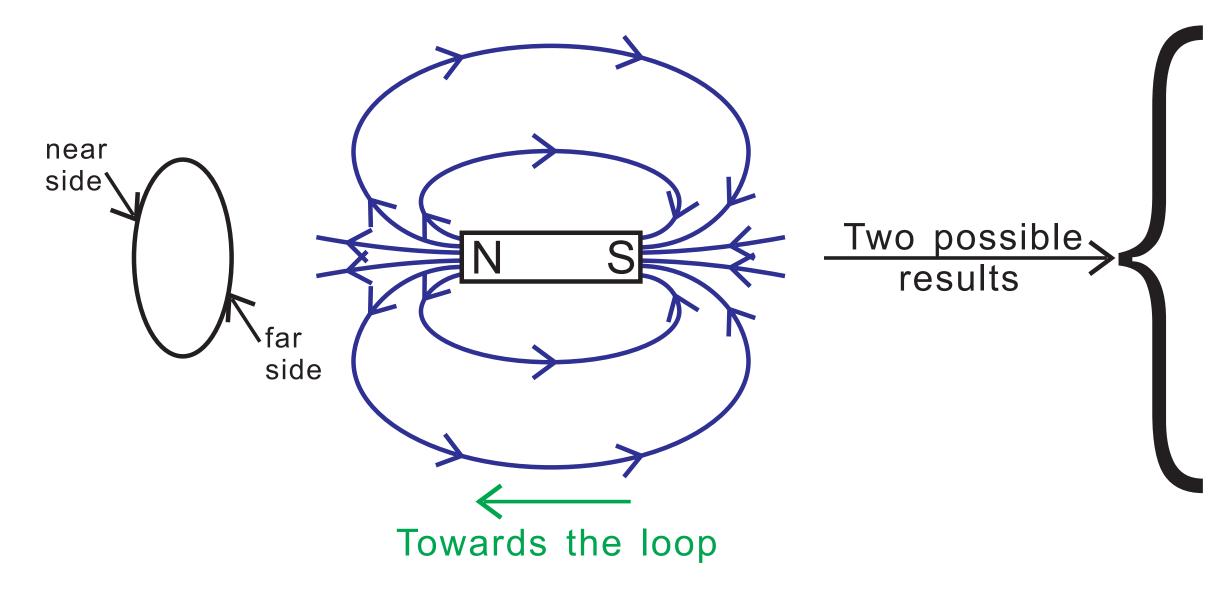
 (\mathbf{d}) loop, a current is induced in the direction shown. The magnetic field lines stormal flux are those due to the bar magnet. (b) This induced current produces its own magnetic field directed to the left that counteracts the increasing external flux. The field lines shown are those due to the induced current in the ring. ha magnet is moved away from the stationary conducting loop, a current n the direction shown. The magnetic field lines shown are those due to net $\frac{1}{d}$ This induced current produces a magnetic field directed to the counteraques the decreasing external flux. The magnetic field lines hose due to the induced current in the ring.

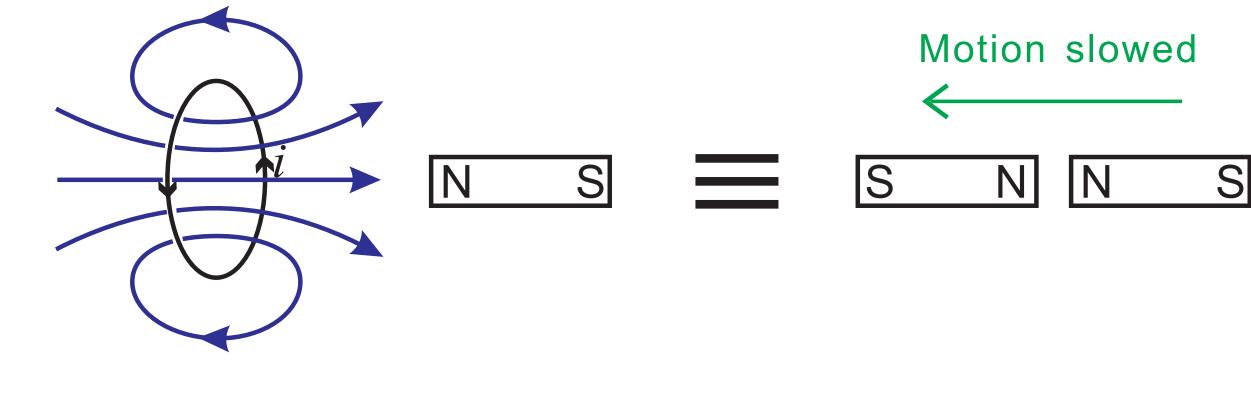
> This induced current produces magnetic field directed to the right and so counteracts decreasing ed lines shown are those due to induced current in ring

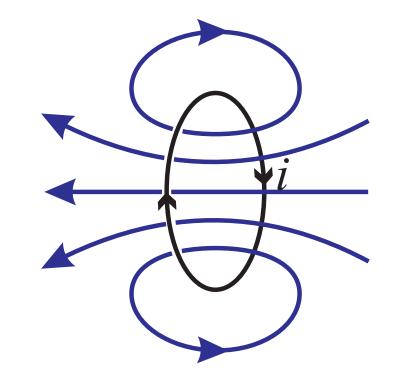




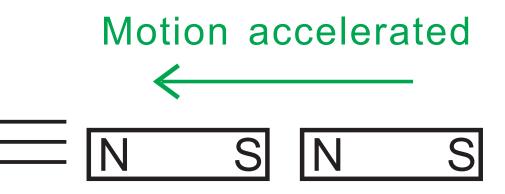




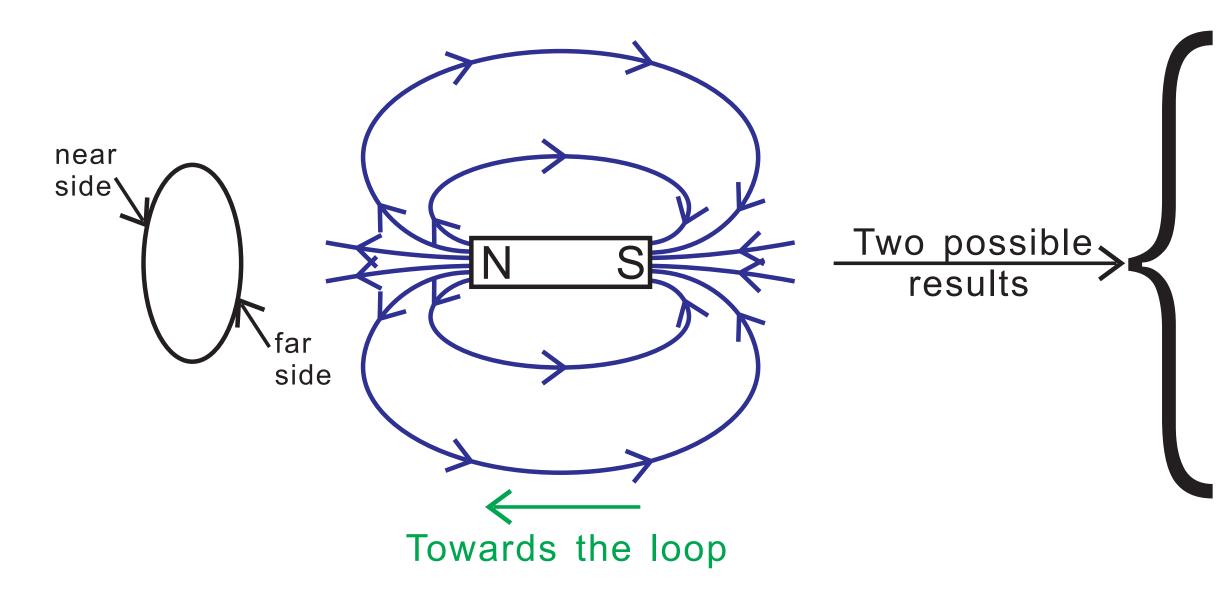


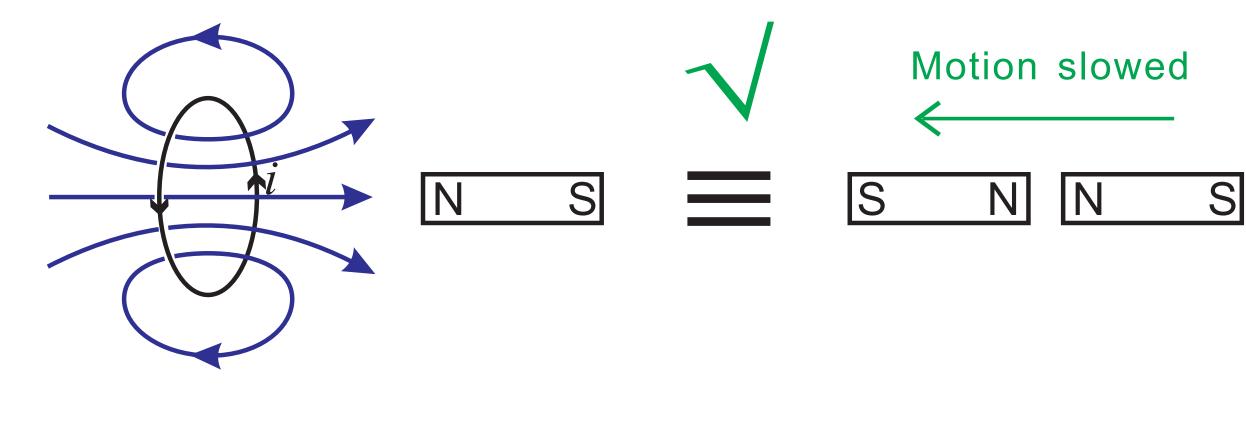


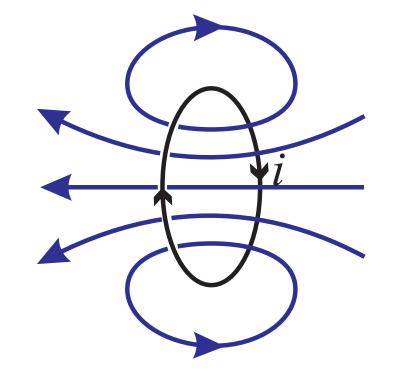




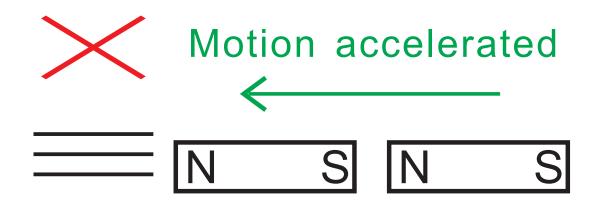






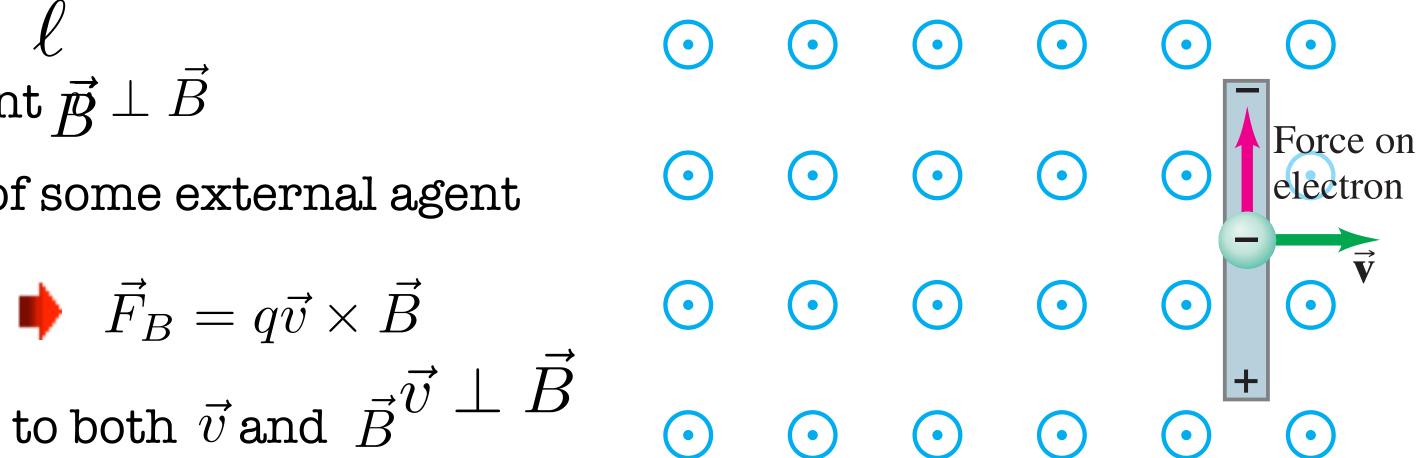






Motional ENF

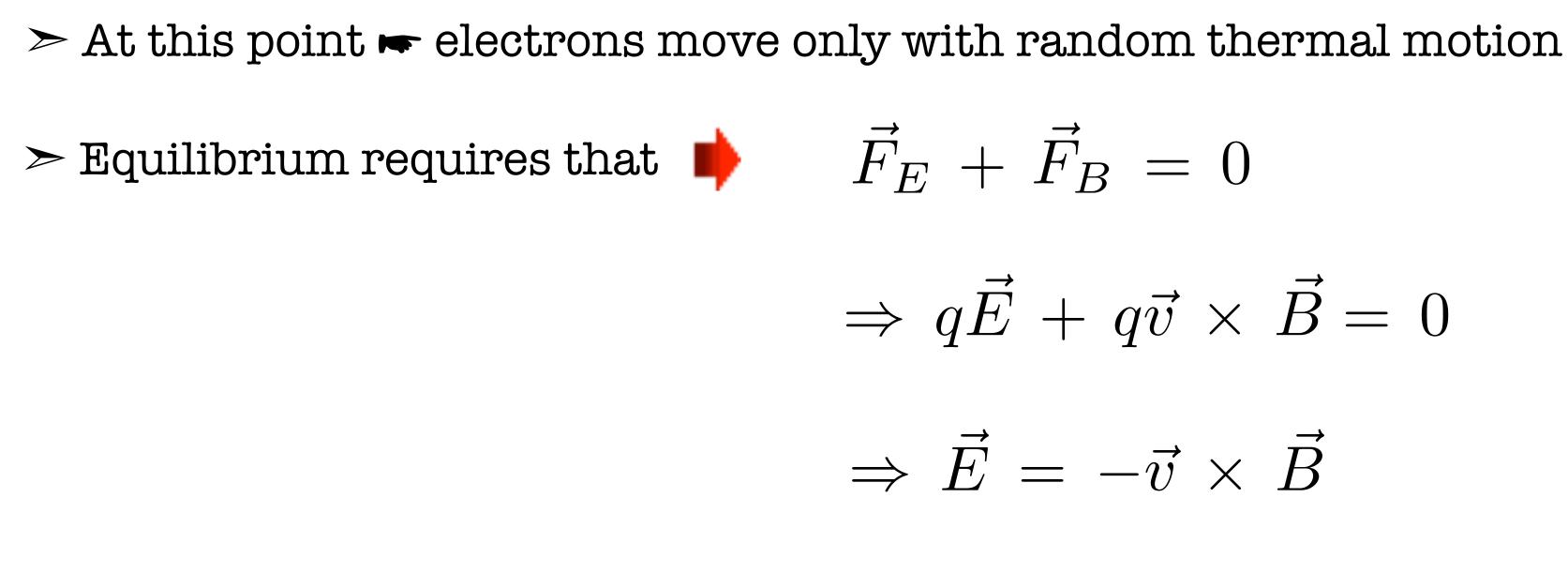
- > Straight conductor of length ℓ is moving through uniform \dot{B} -field directed out of the page > Assume conductor is moving with constant $\vec{R} \perp \vec{B}$
 - under influence of some external agent
- > Electrons in conductor experience force
 - directed along the length perpendicular to both $\vec{v}\, {\rm and}\,\, \vec{B}^{\vec{v}} \perp \vec{B}$
- > Under influence of this force electrons move to upper end of conductor and accumulate there leaving net positive charge at lower end
- \succ Because of this charge separation electric field E
- is balanced by the upward electric force qE



$$ec{F}_B = qec{v} imes ec{B}$$
 is produced inside conductor $ec{v} \ ec{B}$ $ec{B}$

 \succ Charges accumulate at both ends until downward magnetic force qvB on charges remaining in conductor





> Voltage across ends of conductor \blacktriangleright $\Delta V = -E\ell$

 \therefore Voltage $\blacktriangleright \mathcal{E} = \Delta V = vB\ell$

Potential difference is maintained between ends of conductor as long as the conductor continues to move through the uniform magnetic field

 $\Rightarrow q\vec{E} + q\vec{v} \times \vec{B} = 0$

 $\Rightarrow \vec{E} = -\vec{v} \times \vec{B}$



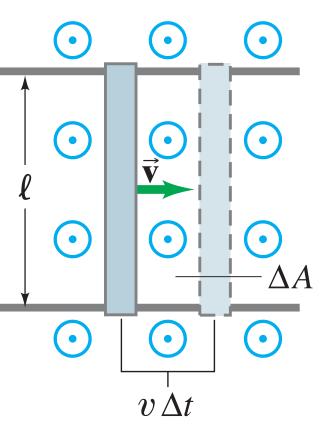
ENF Induced in Moving Conductor

- and movable road resting on it (\cdot) \odot (\cdot) (\bullet) \odot (\cdot) (\cdot) (\cdot) (\cdot) (\cdot) **B** (outward)

$$\mathcal{E} = \frac{\Delta \mathbf{r}_{M}}{\Delta t} = \frac{D \Delta T}{\Delta t}$$

> If rod is made to move at speed v to right right in travels $\Delta x = u \Delta t$ in time Δt $\Delta x = v \Delta t$ > Area of loop increases - $A = \ell \Delta x = \ell v \Delta t$ in time $\Delta t = \ell \Delta x = \ell v \Delta t$ Δt \mathcal{E} > By Faraday's law $rac{r}$ there is induced emf \mathcal{E} whose magnitude is $R \Lambda A$ $\Lambda \Phi_{\Lambda I}$ $B\ell v\Delta t$ $= B\ell v$ Δt > Induced current is clockwise \Rightarrow (to counter the increasing flux)

> Assume that uniform magnetic field B is perpendicular to area bounded by U-shaped conductor



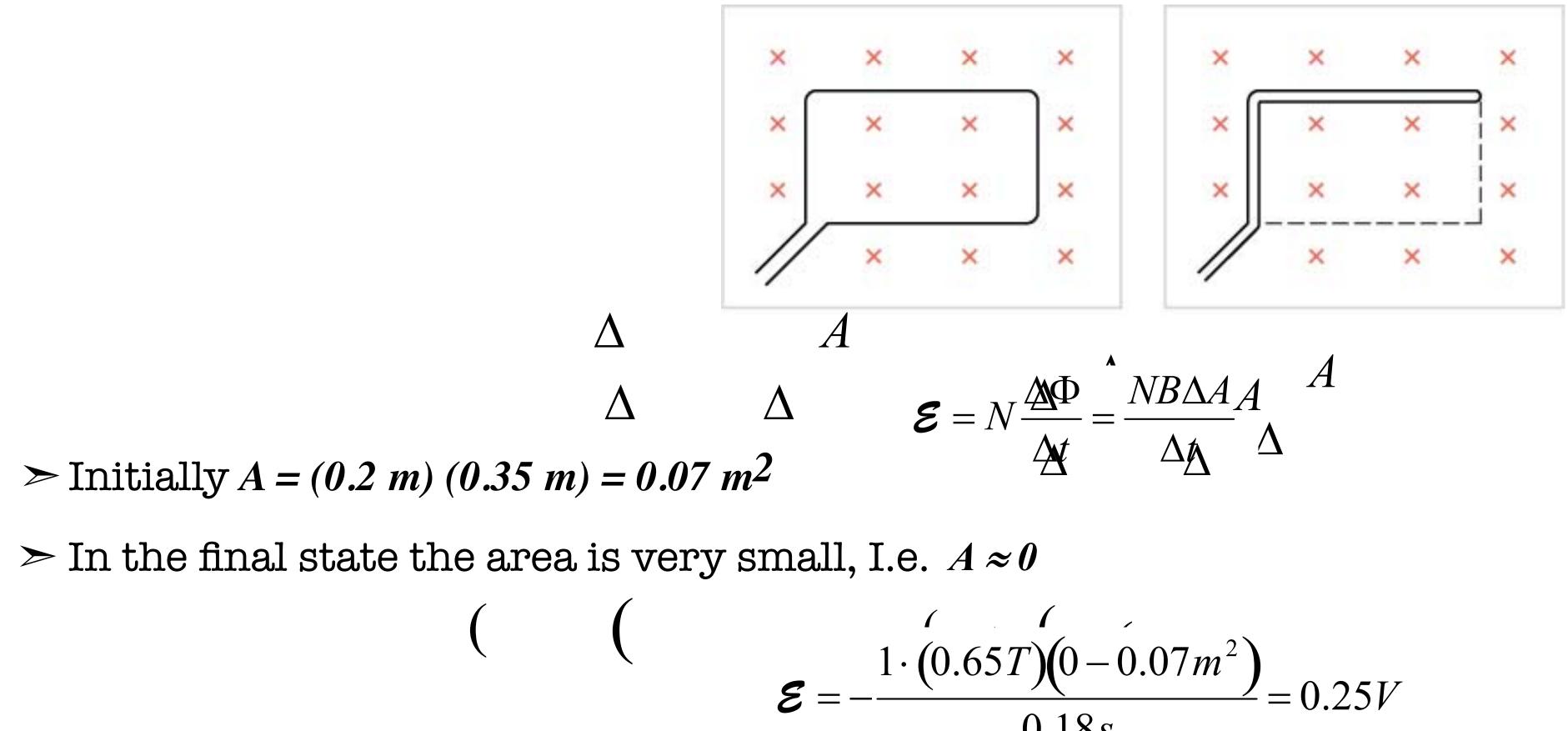




> A rectangular loop of wire with sides of 0.20 and 0.35 m lies in a plane perpendicular to a constant magnetic field of 0.65 T

> In a time of 0.18 s, one-half of the loop is folded back on the other half as shown

> What is the average EMF induced in the loop?



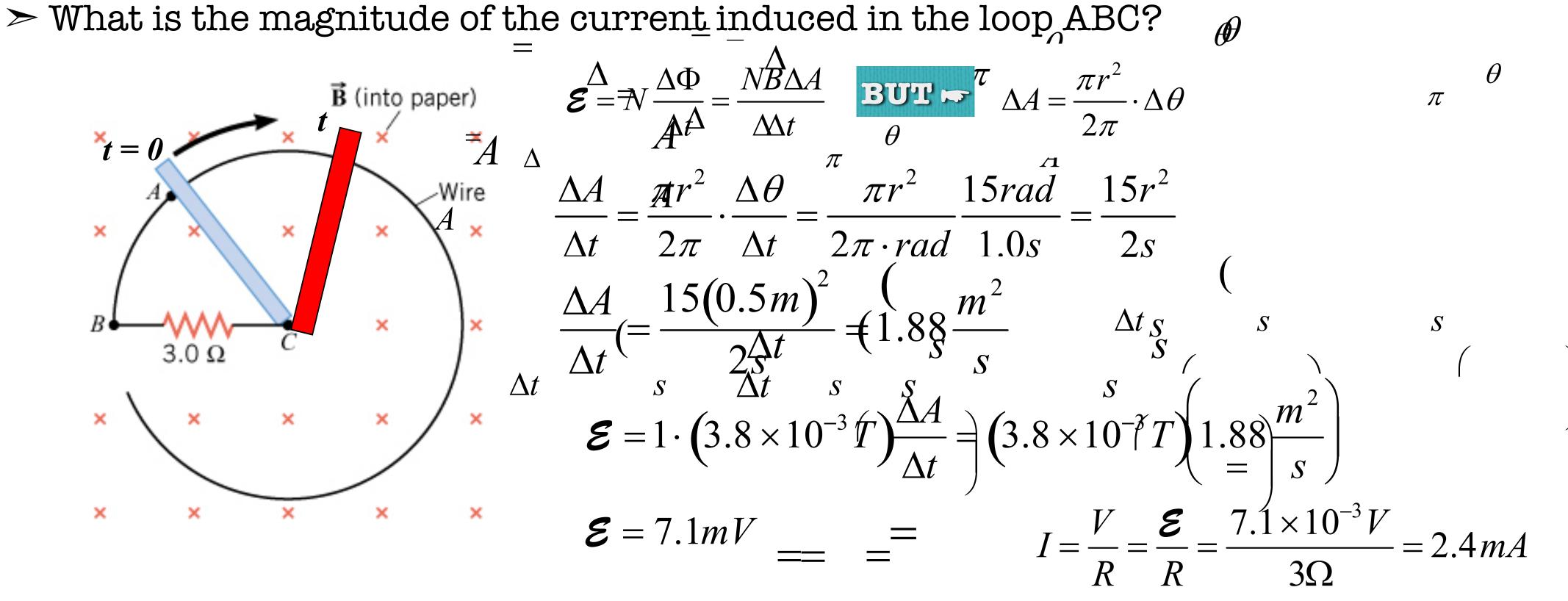
$$N \frac{\Delta \Phi}{\Delta t} = \frac{NB\Delta A}{\Delta t} \Delta A$$

$$\frac{(55T)(0-0.07m^2)}{0.18s} = 0.25V$$



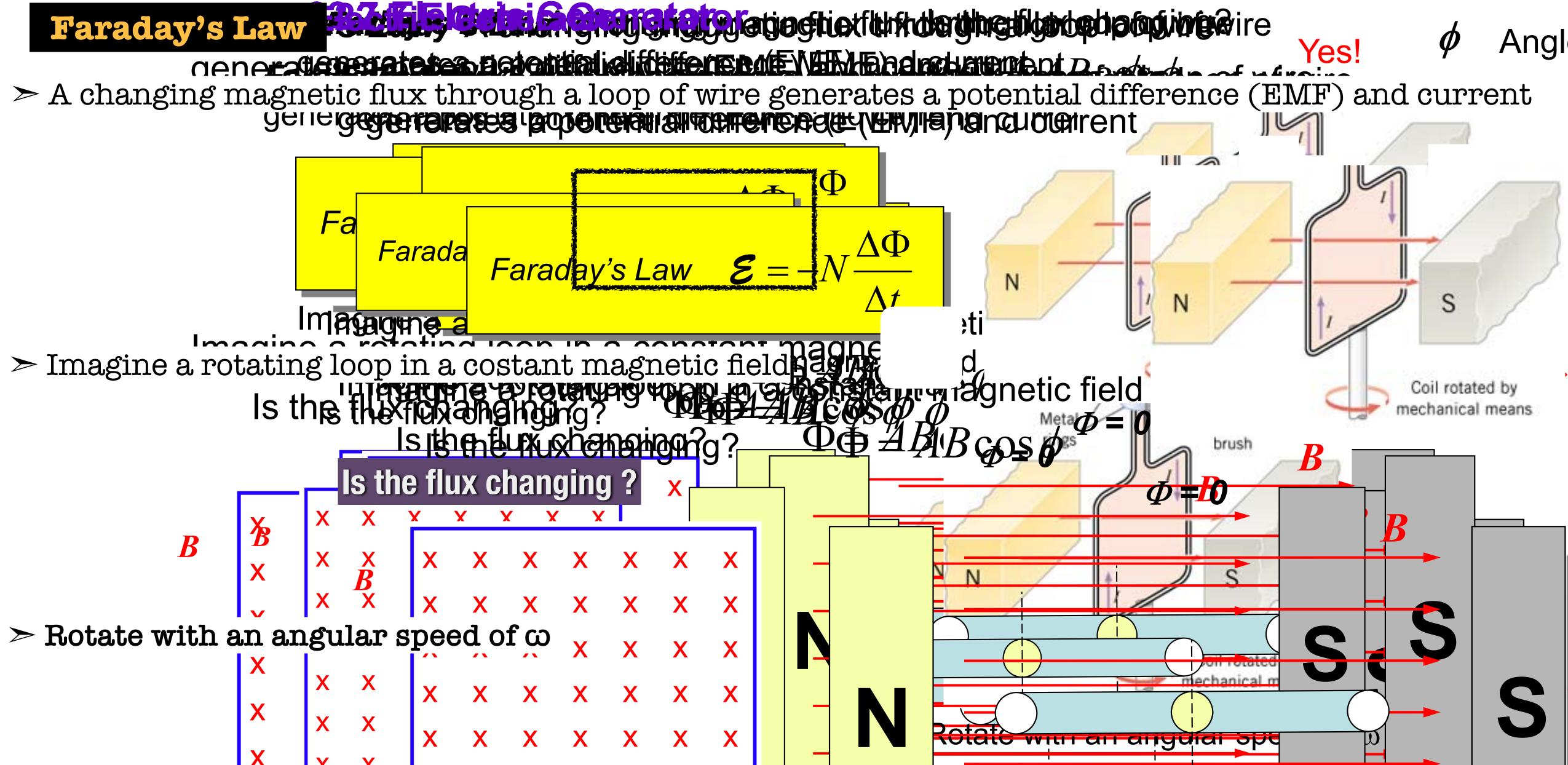
Example

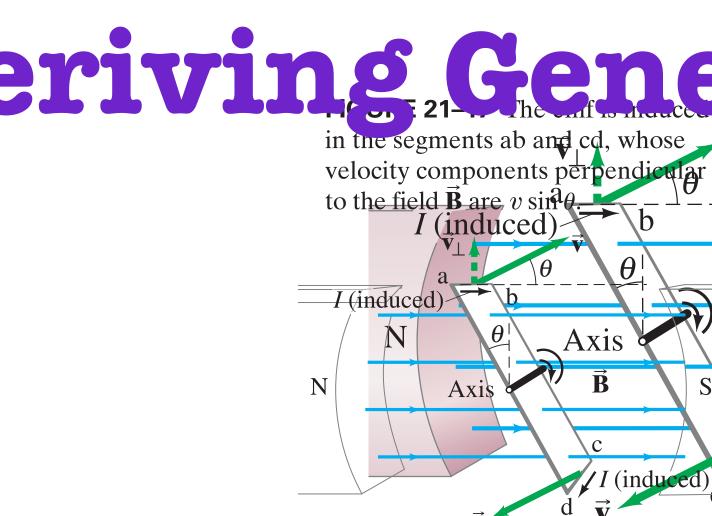
- > The drawing shows a copper wire bent into a circular shape with a radius of 0.5 m > Radial section BC is fixed in place, while copper bar AC sweeps around at an angular speed of 15 rad/s > The bar makes electrical contact with the wire at all times > The wire and the bar have negligible resistance > A uniform magnetic field exists everywhere, is perpendicular to plane of circle and has a magnitude of 3.8 x 10-3 T





22.7 Etectric Generator





not along wire's length \Rightarrow emf generated is due only to force on charges in sections ab and cd

> RHR - direction of the induced current in ab is from a toward b

- > In the lower section \blacktriangleright it is from c to d \Rightarrow flow is continuous in the loop $v_{\perp} = v \sin \theta \Rightarrow \mathcal{E} = 2NB\ell \sin \theta \quad \mathcal{E} = B\ell v_{\perp}$ > Magnitude of the emf generated in ab is $\blacktriangleright \mathcal{E} = B\ell v_{\perp}$
- emf induced in cd has same magnitude and is in same direction 🖛 emfs add

 $v_{\perp} = v \sin \theta \Rightarrow v \mathcal{E} \equiv 2 N \mathcal{B} \ell \vec{v} \sin \theta \sin \theta$

anemating current in the stator cons, which is the output. This ac output is changed to dc for charging the battery by the use of semiconductor diodes, which allow current flow in one direction only.

Deriving Generation of the second of the with the with the period of the second of the

made to rotate clockwise in a uniform magnetic field \vec{B} . The velocity of the two lengths ab and cd at this instant are shown. Although the sections of wire bc and da are moving, the force on electrons in these sections is toward the side of the wire, not along the wire's length. The emf generated is thus due only to the force on charges in the sections ab and cd. From right-hand-rule-3, we see that the direction of the induced current in ab is from a toward b. And in the lower section, It is from c so the flow is continuous in the loop. The magnitude of the emf generated in ab is given by Eq. 21–3, except that we must take the component of the velocity perpendicular to *B*:

$\mathcal{E} = B\ell v_{\perp},$

where ℓ is the length of ab. From Fig. 21–17 we see that $v_{\perp} = v \sin \theta$, where θ is the angle the loop's face makes with the vertical. The emf induced in cd has the same > Loop is being made to rotate clockwise in a uniform magnetic field. Therefore their emfs add, and the total emf is

where we have multiplied by N, the number of loops in the coil.

- > Velocity of two lengths ab and cd at this instant are showing is rotating with constant angular velocity ω , then the angle $\theta = \omega t$. From the angular equations (Eq. 8-4), $v = \omega r = \omega (h/2)$, where r is the distance from the rotation axis and h is the length of bc or ad. Thus $\mathscr{E} = 2NB\omega\ell(h/2)\sin\omega t$, or
- > Although sections of wire bc and da are moving reference of e's in these sections is toward side of wire where $A = \ell h$ is the area of the loop. This equation holds for any shape coil, not just where $A = \ell h$ is the area of the loop. This equation holds for any shape coil, not just

- **umber of loops in coil**

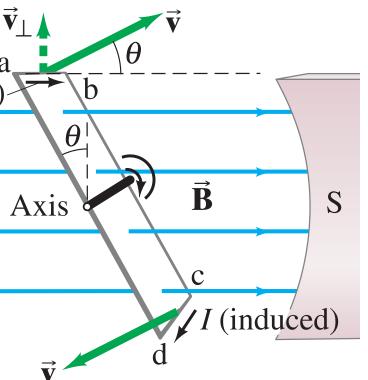


If coil is rotating with constant angular velocity

- $v = \omega r$ =
- r relation distance from the rotation axis hlength of bc or ad

 $A = \ell h$ rea of loop

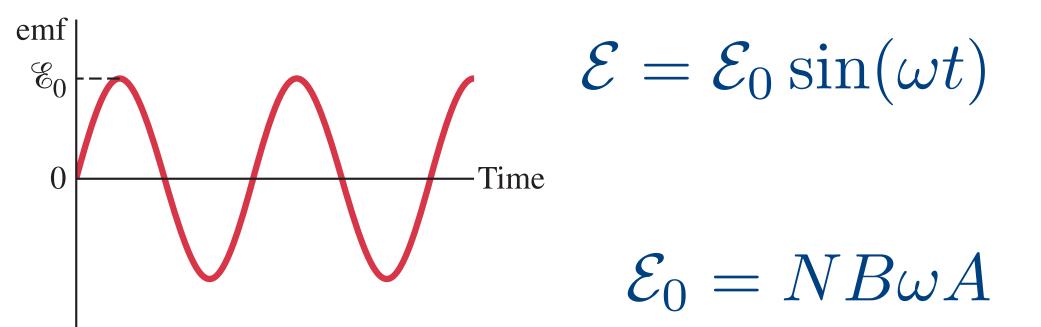
$$V_{\rm rms} = \frac{NB\omega A}{\sqrt{2}}$$

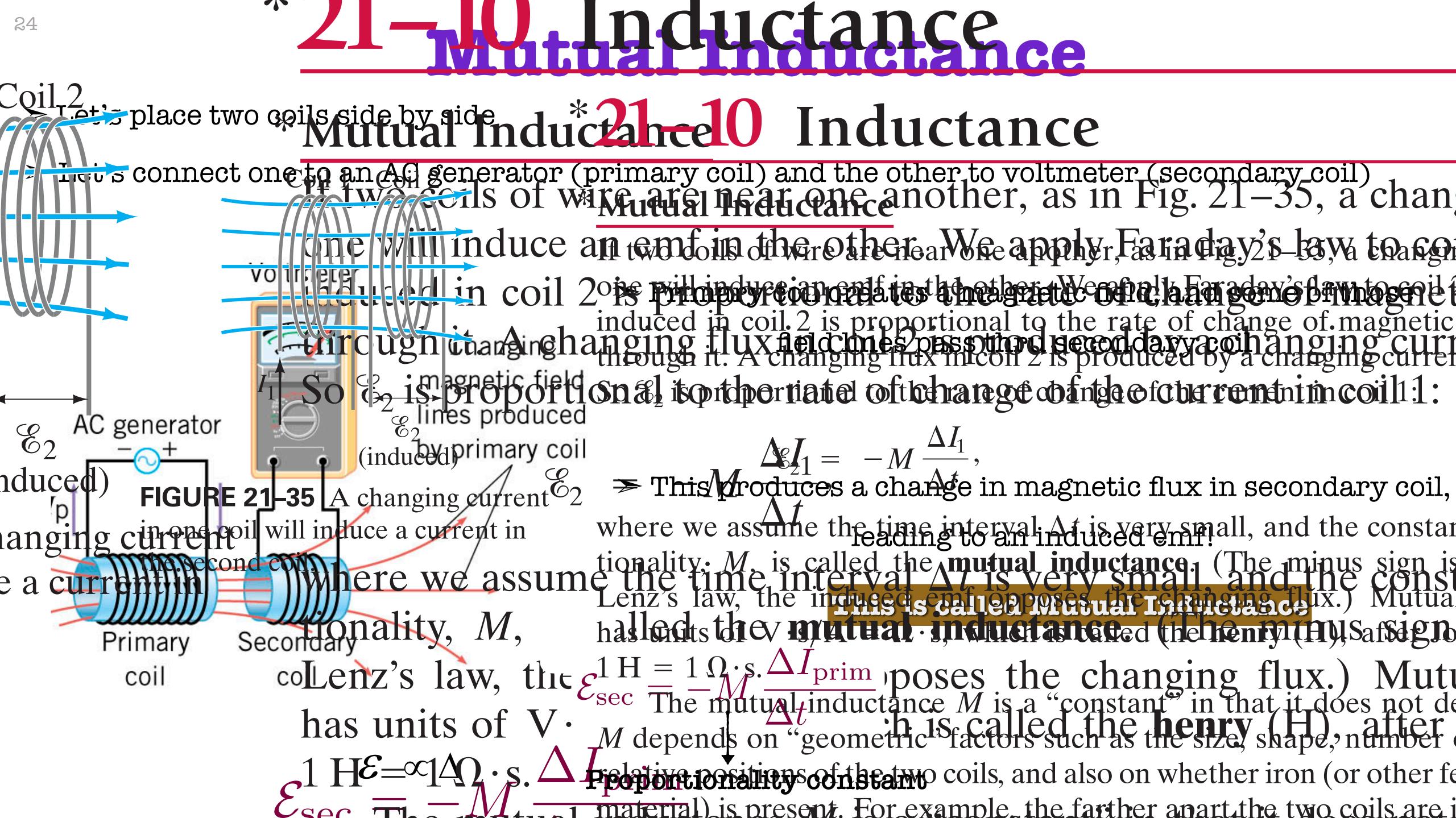


y
$$\omega \Rightarrow \theta = \omega t$$

= $\omega(h/2)$

$\mathcal{E} = 2NB\omega\ell(h/2)\sin(\omega t) \Rightarrow \mathcal{E} = NB\omega A\sin(\omega t)$





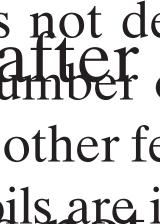
21-hugggese

the will induce an temfoin the other apply, Faraday's by toago Fillin coil 2019 priviperreiznordates anedgratic oppelates anedgratic oppelates and a serve barragene induced in coil 2 is proportional to the rate of change of magnetic ugn changing hanging flux field interpresented avy coin anoing curr So is propositional itomate of the second states of the second states of the second se

 $\sum_{i=1}^{\infty} I_{i} = -M \frac{\Delta I_{1}}{\Delta f},$ This produces a charge in magnetic flux in secondary coil, where we assume the time interval At is very small, and the constant collenz's law, the $\mathcal{E}_{sec}^{1H} = 10$ s. ΔI_{prim} poses the changing flux.) Muture has units of $V \cdot M$ depends on "geometric" factors such as the size, shape, number $\mathcal{E}_{sec}^{1H} = 10$ s. $\Delta P_{sec}^{1H} = 10$ s. ΔI_{prim} poses the changing flux.) Muture has units of $V \cdot M$ depends on "geometric" factors such as the size, shape, number $\mathcal{E}_{sec}^{1H} = 100$ s. $\Delta P_{sec}^{1H} = 100$ s







²⁵ the advanta Serv P replace a battery when ≻ Consider just one coil connected to an AC gener
> AC current produces afchanging magnetic field
The concept of inducta
Ac generator Ac generator Ac generator Ac generator Ac generator Ac generator Ac generator Ac generator Ac generator Inside the coil, and the change coin flux (Lenz's (For example, if the control of the co
$\mathcal{E} = -L \frac{\Delta I}{\Delta t} \int_{C}^{C} \mathbf{P} \mathbf{P} \mathbf{P} \mathbf{P} \mathbf{P} \mathbf{P} \mathbf{P} P$

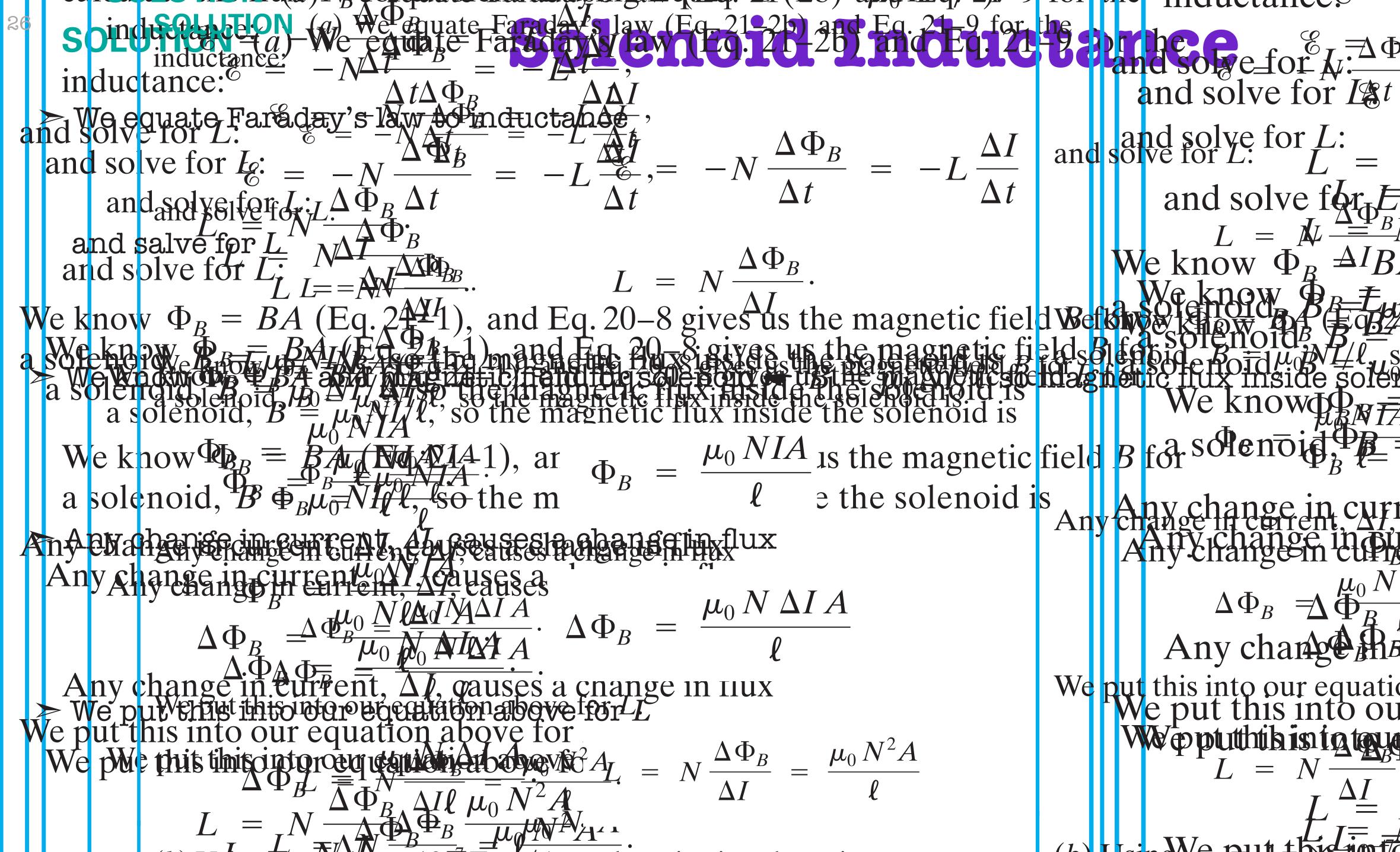
mutual muutanee to a second con m the pacemaker near the ne starrage makers in that surgery n it wears out. erador

d which produces a change in magnetic flux within coil

ance applies also to an isolated single coil. the solution of the solution o 's law): it is much like the back emf gener This process is called Self Induction current through the coil is increasing, the emf that opposes the original current an uced emf & is proportional to the rate of o n opposed to the change, hence the minus rtionality constant

The constant of proportionality L is called the self-inductant



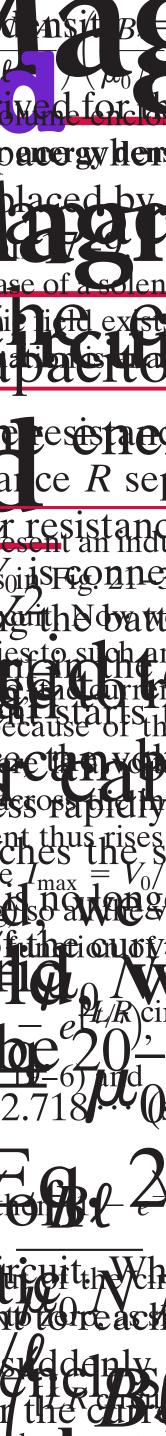


Soluction Eq. (a) We equate Farada is law (Eq. 21-2b) and Eq. 21-9 for the participant of the part and solve $f \phi_T E$: $M \phi_R$ We know $\Phi_R \stackrel{\frown}{=} \Delta IBA / \Delta Ia$. We know $\Phi_B = BA$ (Eq. 241), and Eq. 20-8 gives us the magnetic field we know $B_B = E_B BA$ (Eq. 241), and Eq. 20-8 gives us the magnetic field we know $B_B = E_B BA$ (Eq. 241). We know $\overline{A} \overline{A}$ e magnetic flux inside the solenoid is We know $\Phi_B = BA(NIA 1)$, at $\Phi_B = \frac{\mu_0 NIA}{\ell}$ is the magnetic field *B* for solenoid $B_B = \frac{\mu_0 NIA}{\ell}$ a solenoid, $B \Phi_B \mu_0 NIA$, so the m $\Phi_B = \frac{\mu_0 NIA}{\ell}$ is the solenoid is Any change in current AAny change in thent, MI, $\Delta \Phi_B = \Delta \Phi_B = \mu_0 N \Delta I A \mu_0 I$ $M choose \Phi B = \mu_0 N \Delta I A \mu_0 I$ Any change Reutrent We put this into our equation above We put this into our equat Weputhibisistateupageatin $\mu_0 N^2 A$



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Energy density = $\frac{1}{2}\varepsilon_0 E^2 + \frac{1}{2\mu_0}B^2$

Induced Electric Field×

So far we have discussed that a *change* in magnetic flux will lead in an induced emf distributed in the loop, resulting from an induced E-field. > We have seen that a changing magnetic flux

induces an emf and a current in a conducting loop

 $\therefore B \uparrow \Rightarrow$ anti-clockwise current However, even in the absence of the loop (so that the societation is included current), the induced E-field will still accompany a change in magnetic flux. > In the same way we want felate in the conducting loop 105

to an electric field by claiming that electric field is created in

conductor as a result of the changing magnetic flux

 \therefore Consider a circular path in a region \succ $\mathcal{E} = \sum E_{\parallel} \oplus E_{\parallel} \oplus E_{\parallel} \oplus E_{\parallel} \oplus E_{\parallel}$ Non-conservative force field $\begin{array}{c} closed \\ path \end{array}$

