# School of Physics 2013 NSF-CMACS Workshop on Atrial Fibrillation

Georgia

Tech

CMACS

Computational Modeling and Analysis for Complex Systems

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# Outline

- Motivation
- Overview of Excitable Systems
- Cardiac dynamics as an excitable system
- Chaos and tools from nonlinear dynamics can be used to study cardiac arrhythmias.
- Modeling cardiac dynamics
- Initiation, dynamics and instabilities of spiral waves
- Defibrillation and new methods to control arrhythmias using control methods from Chaos theory.
- Get you ready for your projects



Heart disease is one of the leading causes of death in the world.

Ranks number one in industrialized countries.

In the USA alone:

- 1 in 5 have some form of heart disease.
- 4.5 million do not die but are hospitalized every year.
- Economic impact: \$214 billion a year.
- 1/3 of total deaths are due to heart disease.

As computer science, Mathematics & Statistics and Biology majors : This is a great excitable system with a broad range of dynamical states where you can apply your area of expertise and have an impact, as there are countless of remaining exciting open questions .

# Nonlinear Dynamics: Complex Systems and Pattern Formation



Examples









# Nonlinear Dynamics: Complex Systems and Pattern Formation







Suzuki N et al. PNAS 2003; 100:9680-9685













Excitable media can be defined as a system formed by segments or elemental cells, each of which has the following properties:

- A well defined steady state (or rest state)
- A threshold for excitation
- An excitation that only a function of the system and independent of the stimulus.
- A refractory period following excitation
- A diffusive coupling with near neighbors ( in space )

They can be: Physical, biological, mechanical, electrical, mathematical, chemical, ... and they can lead to chaos.

#### Example: Mexican Wave







#### Example: Forest fires







Good example to understand refractory period and collision between waves. (not solitons)

Problem: very long wave length, even longer refractory period

Example: Toilet



Probably best example: Threshold,

Excitation (that is independent of the stimulus) Refractory period

Extras: complex dynamics such as alternans.

Problems: Only good for 0D, difficult to couple in higher dimensions.

#### Example: Predator-Pray (oscillators)



Oscillators can be a subclass of Excitable media.

Very useful to understand qualitatively their dynamics



Huffaker (1958) reared two species of mites to demonstrate these coupled oscillations of predator and prey densities in the laboratory. Using *Typhlodromus occidentalis* as the predator and the six-spotted mite (*Eotetranychus sexmaculatus*) as the prey

#### Example: Saline oscillator (also density oscillator)



Oscillators can behave as excitable systems when externally excited at frequencies higher than their Natural oscillation.



100

200

300

### **Complex Systems and Pattern Formation**

One of the simplest system with complex dynamics (and the cheapest!)



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#### the Saline Oscillator





### **Complex Systems and Pattern Formation**

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# We get to play with them today!

#### Example: Chemical oscillators





Many types of chemical reactions that oscillate and have interesting dynamics in space.

We get to play with them tomorrow!

#### Example: excitable biological cells





Neurons

Myocites (cardiac cells)

In these systems, the excitability is a change in voltage

We get to study them the rest of the workshop!

#### **Excitable Biological Cells**

In these cells, the voltage is a consequence of a gradient in the cellular membrane due to a difference in concentrations of ions inside and out-side of the cell.





### **Excitable Biological Cells**

The difference in electric potential is a consequence of the intracellular

and extracellular ion concentrations.



#### **Excitable Biological Cells**

The difference in electric potential is a consequence of the intracellular

and extracellular ion concentrations.



LINK to applet

Resting membrane potential (Steady State)

### Electrical Activity in Cells

How does electricity play a role?

- The cell membrane: lipid bi-layer 10 nm thick, impermeable to ions except through specialized proteins (ion channels).
- Ion concentration gradient and voltage drop across membrane.
- Movement of ions across the membranes produces an action potential.
- Active transport through pumps and exchangers in the membrane restores original concentrations.





### Electrical Activity in Neurons

Na⁺, K⁺ , Cl⁻,







### Electrical Activity in Neurons

#### Transfer of information is the main consequence of electrical activity in neurons.





How does electricity play a role in the heart?

 Cardiac cells are about 100-150 µm in length, 10-20 µm in diameter.



#### Ca<sup>2+</sup>, Na<sup>+</sup>, K<sup>+</sup>







Ca<sup>2+</sup>, Na<sup>+</sup>, K<sup>+</sup>









#### **Electrical-Contraction Coupling**

- Cellular action potential triggers contraction through calcium processes.
- Increased calcium current stimulates release of intracellular store.
- Transiently increased calcium binds to contraction proteins.





Contraction of the heart occurs because the electrical AP propagate as waves thought the heart.

### **Electrical Waves in Tissue**

# Cells connected in a 2D preparation





$$i_{intracell}(x-\Delta x,t)\Delta x \qquad i_{intracell}(x,t)\Delta x \qquad i_{intracell}(x+\Delta x,t)\Delta x$$

$$f_{g}(x-\Delta x/2)\Delta x \qquad f_{g}(x-\Delta x/2)\Delta x \qquad r_{g}(x+\Delta x/2)\Delta x \qquad r_{g}(x+\Delta x/2)\Delta x$$

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$$f_{g}(x-\Delta x,t)\Delta x \qquad f_{g}(x+\Delta x,t)\Delta x \qquad f_{g}(x+\Delta x,t)\Delta x$$

 $i_{extracell}(x-\Delta x,t)\Delta x$   $i_{extracell}(x,t)\Delta x$ 

Sum of the ionic currents that make up the action potential.

 $i_{extracell}(x+\Delta x,t)\Delta x$ 

Tissue connection and structure information is given by the diffusion coefficient D.

ion

#### The first mathematical model of electrical AP

J. Physiol. (1952) 117, 500-544

#### A QUANTITATIVE DESCRIPTION OF MEMBRANE CURRENT AND ITS APPLICATION TO CONDUCTION AND EXCITATION IN NERVE

BY A. L. HODGKIN AND A. F. HUXLEY

From the Physiological Laboratory, University of Cambridge

(Received 10 March 1952)

This article concludes a series of papers concerned with the flow of electric current through the surface membrane of a giant nerve fibre (Hodgkin, Huxley & Katz, 1952; Hodgkin & Huxley, 1952 a-c). Its general object is to discuss the results of the preceding papers (Part I), to put them into mathematical form (Part II) and to show that they will account for conduction and excitation in quantitative terms (Part III).

PART I. DISCUSSION OF EXPERIMENTAL RESULTS

J. Physiol. Vol 117 500-540, 1952 by A.L. Hodgkin and A.F. Huxley.

Ca<sup>2+</sup>, Na<sup>+</sup>, K<sup>+</sup>









#### How to model the Neuron AP?

#### The Hodgkin-Huxley model of four variables for neurons



Capacitance is a measure of the amount of electrical energy stored (or separated) for a given electric potential, where C = Q/V

 $Q_m = CV_m$ ;  $dQ_m/dt = I_ion$ ; Therefore  $I_ion = C dV_m/dt$ 

 $dV_m/dt = I_ion/C$ 

J. Physiol. Vol 117 500-540, 1952 by A.L. Hodgkin and A.F. Huxley.

#### How to model the Neuron AP mathematically?

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$$egin{aligned} &I_K &= n^4 \mathbf{g_k} \left( V_m - V_K 
ight) \ &I_{Na} &= m^3 h \mathbf{g_{Na}} \left( V_m - V_{Na} 
ight) \ &I_{Cl} &= \mathbf{g_{Cl}} \left( V_m - V_{Cl} 
ight) \end{aligned}$$

Cell currents for the model (follow Ohms law)

Ecuaciones para la probabilidad de las puertas

$$dy/dt~=~lpha_y(V_m)~(1-y)~-~eta(V_m)~y$$

$$lpha_y(V_m)$$
  
 $eta_y(V_m)$ 

J. Physiol. Vol 117 500-540, 1952 by A.L. Hodgkin and A.F. Huxley.

 $lpha_m = 0.1(V_m + 35.0)/(1. - e^{(-(V_m + 35.0)/10.0)}) egin{array}{l} eta_m = 4.0 \, e^{(-(V_m + 60.0)/18.0)} \end{array}$ 

$$lpha_h = 0.07 \, e^{(-(V_m + 60.0)/20.0)} \ eta_h = 1./(1 + e^{(-(V_m + 30.0)/10.0)})$$

 $lpha_n = 0.01 (V_m + 50.0) / (1 - e^{(-(V_m + 50.0)/10.0)}) egin{array}{l} eta_n = 0.125 \, e^{(-(V_m + 60.0)/80.0)} \end{array}$ 

# Gating Variables

- Gating variables are time-dependent variables that modify the current conductance.
- Gates vary between 0 and 1; 1 maximizes current and 0 eliminates it.
- Gates follow the following equation:

$$\frac{\mathrm{d}y}{\mathrm{d}t} = \frac{y_{\infty} - y}{\tau_{y}}$$

where  $y_{\infty}(V)$  is the steady-state value and  $\tau_y(V)$  is the time constant for the gate.

$$dy/dt\,=\,lpha_y(V_m)\,(1-y)\,-\,eta(V_m)\,y$$

$$\beta_y(V_m)$$

$$\alpha_y(V_m)$$



$$\tau_i(V) = rac{1}{lpha_i(V) + eta_i(V)}.$$

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#### How to model the Neuron AP mathematically? Hodgkin-Huxley four variable model

S2 30,17.6 17.5 13 IS2 -300, -350 -450 N, gna 70, 40

Change Rest Membrane p. N t=80, gna = 180, 188, 200, 300 How to model the Neuron AP mathematically? Hodgkin-Huxley four variable model

Even with just 4 variables is hard to understand the dynamics

Ideal to decrease the number of variables to the minimum.

How?

m gate is very fast, make it instantaneous function of Voltage

 $n \sim 1-h$  therefore change to an effective variable w = b - h = a n

$$C\frac{\mathrm{d}u}{\mathrm{d}t} = -g_{\mathrm{Na}}[m_0(u)]^3 (b-w) (u-V_{\mathrm{Na}}) - g_{\mathrm{K}} \left(\frac{w}{a}\right)^4 (u-V_{\mathrm{K}}) - g_{\mathrm{L}} (u-V_{\mathrm{L}})$$
$$\frac{\mathrm{d}w}{\mathrm{d}t} = \frac{1}{\tau_w} G(u,w) ,$$

How to model the Neuron AP mathematically? Hodgkin-Huxley four variable model

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$$\frac{\mathrm{d}u}{\mathrm{d}t} = \frac{1}{\tau} \left[ F(u, w) \right]$$
$$\frac{\mathrm{d}w}{\mathrm{d}t} = \frac{1}{\tau_w} G(u, w) \,,$$

#### IMPULSES AND PHYSIOLOGICAL STATES IN THEORETICAL MODELS OF NERVE MEMBRANE

RICHARD FITZHUGH From the National Institutes of Health, Bethesda

ABSTRACT Van der Pol's equation for a relaxation oscillator is generalized by the addition of terms to produce a pair of non-linear differential equations with either a stable singular point or a limit cycle. The resulting "BVP model" has two variables of state, representing excitability and refractoriness, and qualitatively resembles Bonhoeffer's theoretical model for the iron wire model of nerve. This BVP model serves as a simple representative of a class of excitable-oscillatory systems including the Hodgkin-Huxley (HH) model of the squid giant axon. The BVP phase plane can be divided into regions corresponding to the physiological states of nerve fiber (resting, active, refractory, enhanced, depressed, etc.) to form a "physiological state diagram," with the help of which many physiological phenomena can be summarized. A properly chosen projection from the 4-dimensional HH phase space onto a plane produces a similar diagram

$$\label{eq:dudt} \begin{split} \frac{\mathrm{d} u}{\mathrm{d} t} &= \frac{1}{\tau} \left[ F(u,w) \right. \\ & \frac{\mathrm{d} w}{\mathrm{d} t} = \frac{1}{\tau_w} G(u,w) \,, \end{split}$$

$$\frac{\partial V}{\partial t} = \frac{\partial^2 V}{\partial x^2} + (a - V)(V - 1)V - v$$
$$\frac{\partial v}{\partial t} = \epsilon(\beta V - \gamma v - \delta)$$

- FitzHugh in 1960 made various studies of HH in phase space (fixing values of m,n, h)
- Conclusion needed a simpler model to understand the dynamics

He did not do a reduction of HH

Started with Van der Pol relaxation oscillator (1926) and the phase plane model used by Bonhoeffer

 $\ddot{x} + k\dot{x} + x = 0$  Van der Pol added a damping coefficient

$$\dot{x} + c(x^2 - 1)\dot{x} + x = 0$$

$$= \dot{x}/c + x^3/3 - x$$

$$\dot{x} = c(y + x - x^3/3)$$
$$\dot{y} = -x/c$$

у

He used Lineard's transformation

Bonhoeffer -Van der Pol (BVP)  

$$\dot{x} = c(y + x - x^3/3 + z)$$
  
 $\dot{y} = -(x - a + by)/c$ 



$$\frac{\partial \mathbf{V}}{\partial t} = \frac{\partial^2 \mathbf{V}}{\partial x^2} + (a - \mathbf{V})(\mathbf{V} - 1)\mathbf{V} - \mathbf{v}$$
$$\frac{\partial \mathbf{v}}{\partial t} = \epsilon(\beta \mathbf{V} - \gamma \mathbf{v} - \delta)$$



Four floor to ceiling relay racks, With vacuum tubes (that failed around twice a week) And overloaded the air conditining

#### Nullclines and phase space analysis

$$\frac{\partial V}{\partial t} = \frac{\partial^2 V}{\partial x^2} + (a - V)(V - 1)V - v$$
$$\frac{\partial v}{\partial t} = \epsilon(\beta V - \gamma v - \delta)$$



 $\frac{\partial V}{\partial t} = \frac{\partial^2 V}{\partial x^2} + (a - V)(V - 1)V - v$  $\frac{\partial v}{\partial t} = \epsilon(\beta V - \gamma v - \delta)$ Phase space: a .1 .2 .3 .4 .5 Delta .2, 1.5, 1. Eps .01 .02 .03 .04 Eps .01 .005 .002

.001

		S2 10,9,8,7
		a .1 .2 .3 .4 .5 Delta .2,.1.5,1.
		Also by a =1 S2 =0
		Eps .01 .02 .03 .04
		Eps .01 .005 .002 .001 S2 larger
		Eps .0001 T 700 S2 464 463
. Vol 1 445-466, 1961 by R. FitzHugh	Dynamics of FHN model > 2 h class	T2000 S2 600, 900 1000 1500

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